



Analysis of energetic particle driven MHD modes on EAST

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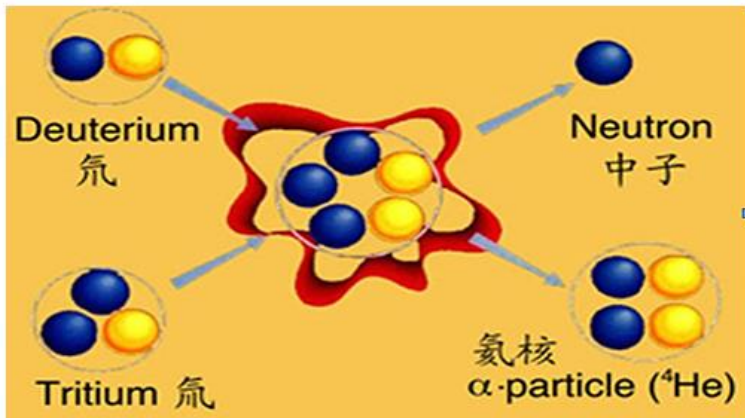
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- ◆ **Motivation**
- ◆ **Kinetic-MHD model in NIMROD**
- ◆ **Functions & progress of AWEAC**
- ◆ **Fishbone/BAE/TAE on EAST**
- ◆ **Conclusions and future works**

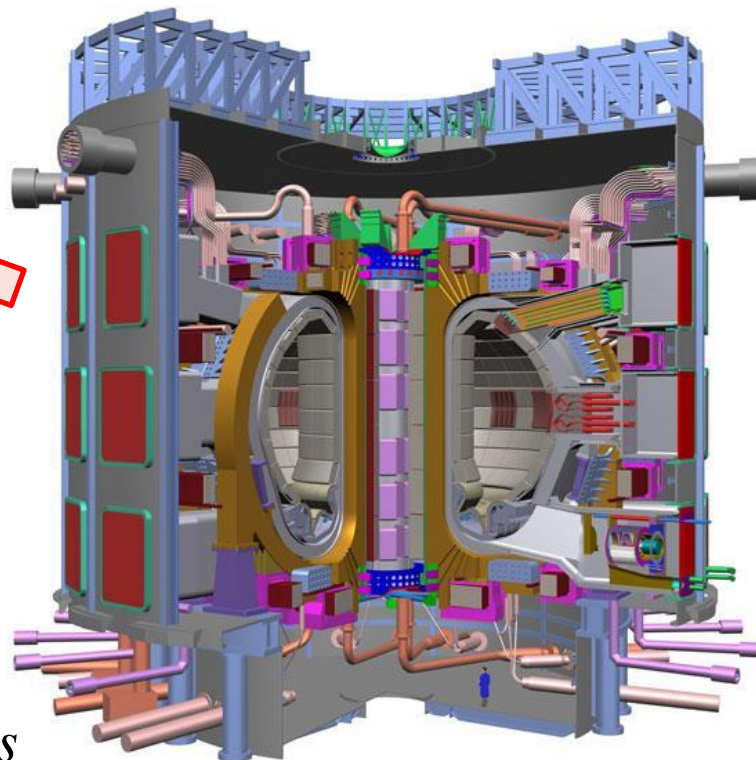
Motivation

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Fusion



ITER



Trapped EP Precession frequency、Core-ion diamagnetic frequency、Passing EP Radial drift velocity/circulating frequency $\sim\sim$

Fishbone frequency/velocity;

EP toroidal velocity $\sim\sim$

AEs phase velocity.

$$0.9 \times 10^6 \text{ m/s} \quad 12 \times 10^6 \text{ m/s}$$

$$\uparrow \quad \uparrow$$

$$V_{Ti} \ll V_A < V_\alpha \ll V_{Te}$$

$$\downarrow \quad \downarrow$$

$$8 \times 10^6 \text{ m/s} \quad 59 \times 10^6 \text{ m/s}$$

$$V_{EP} \sim [3 \times 10^6 \text{ m/s}, 10 \times 10^6 \text{ m/s}]$$

Heating & fusion -> energetic particle

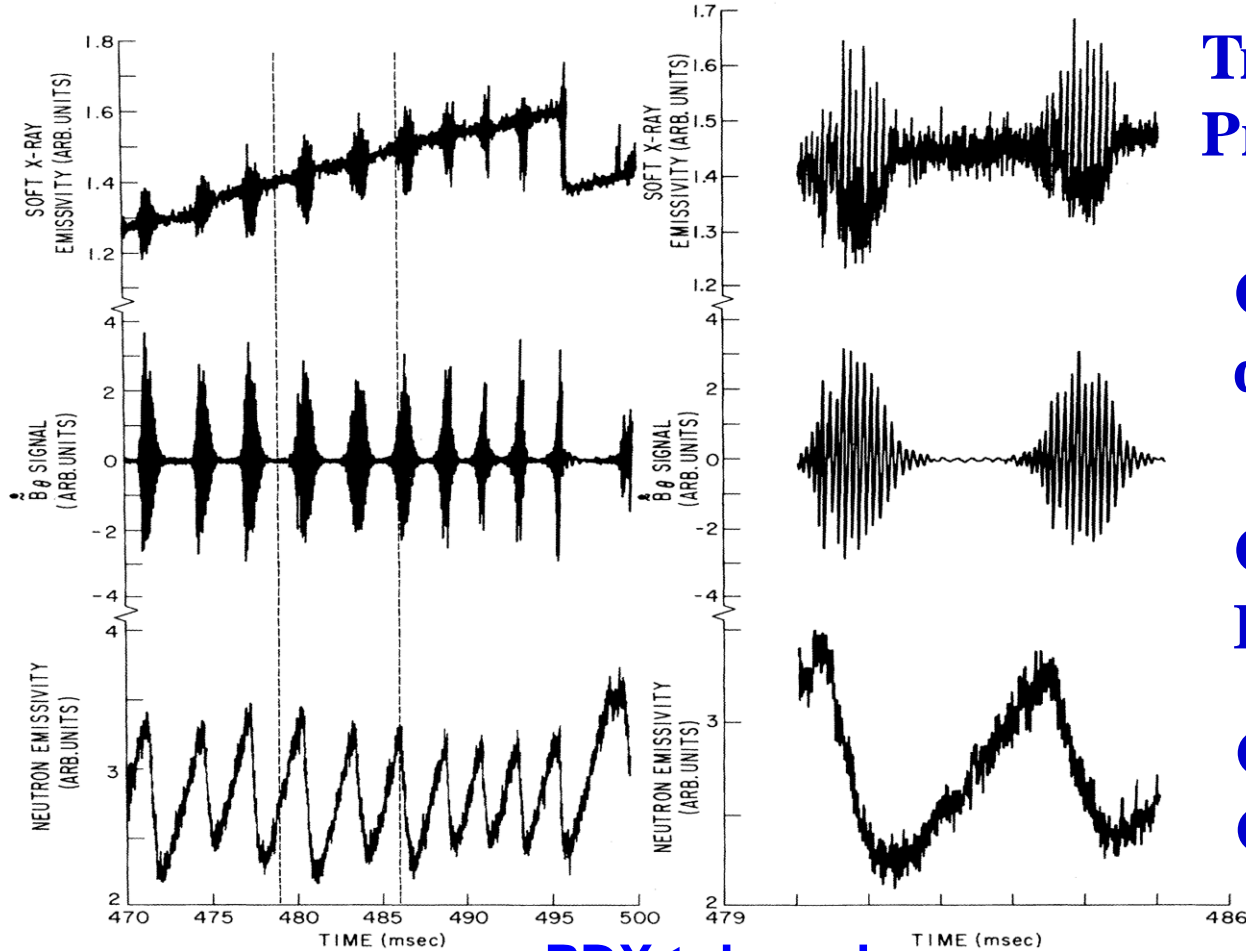
Fishbone -> EP loss $\sim 10\%$
 Alfvén Eigenmodes -> EP loss $\sim 70\%$

1. Heating efficiency -> long pulse steady-state operation;
2. Reduction of EP destruction to FW -> long-lived component;
3. α particle heating -> self-sustained burning plasma.

Fishbone mode

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McGuire et al PRL 1983



PDX tokamak

Trapped energetic ion
Precession frequency

Chen et al PRL 1984

Core-ion
diamagnetic motion

Coppi et al PRL 1986

Circulating energetic ion
Finite radial drift

Betti et al PRL 1993

Circulating energetic ion
Circulating frequency

Wang et al PRL 2001

Superthermal electron

Wong et al PRL 2000

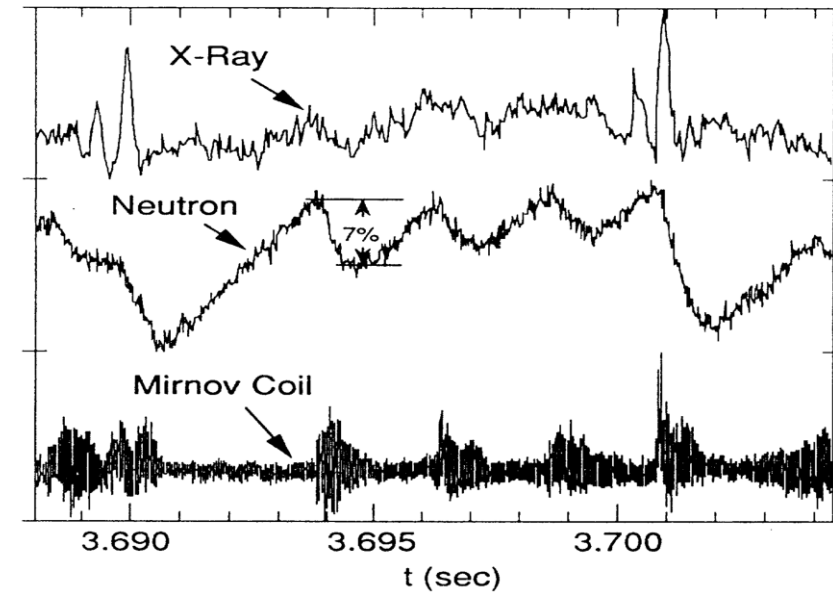
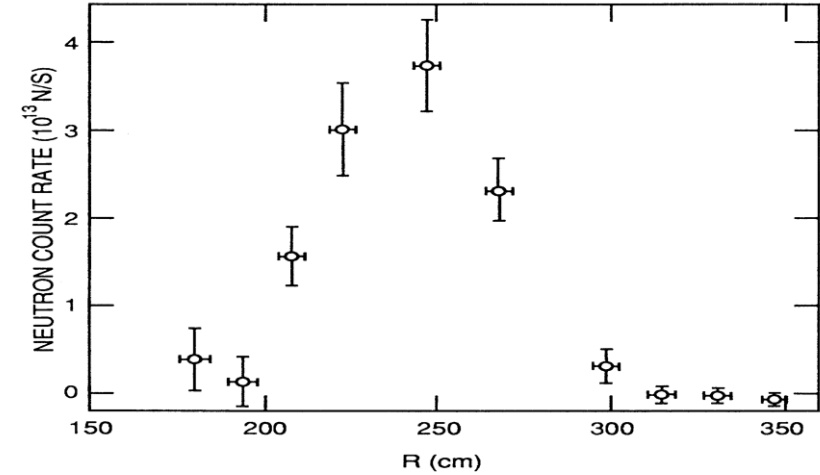
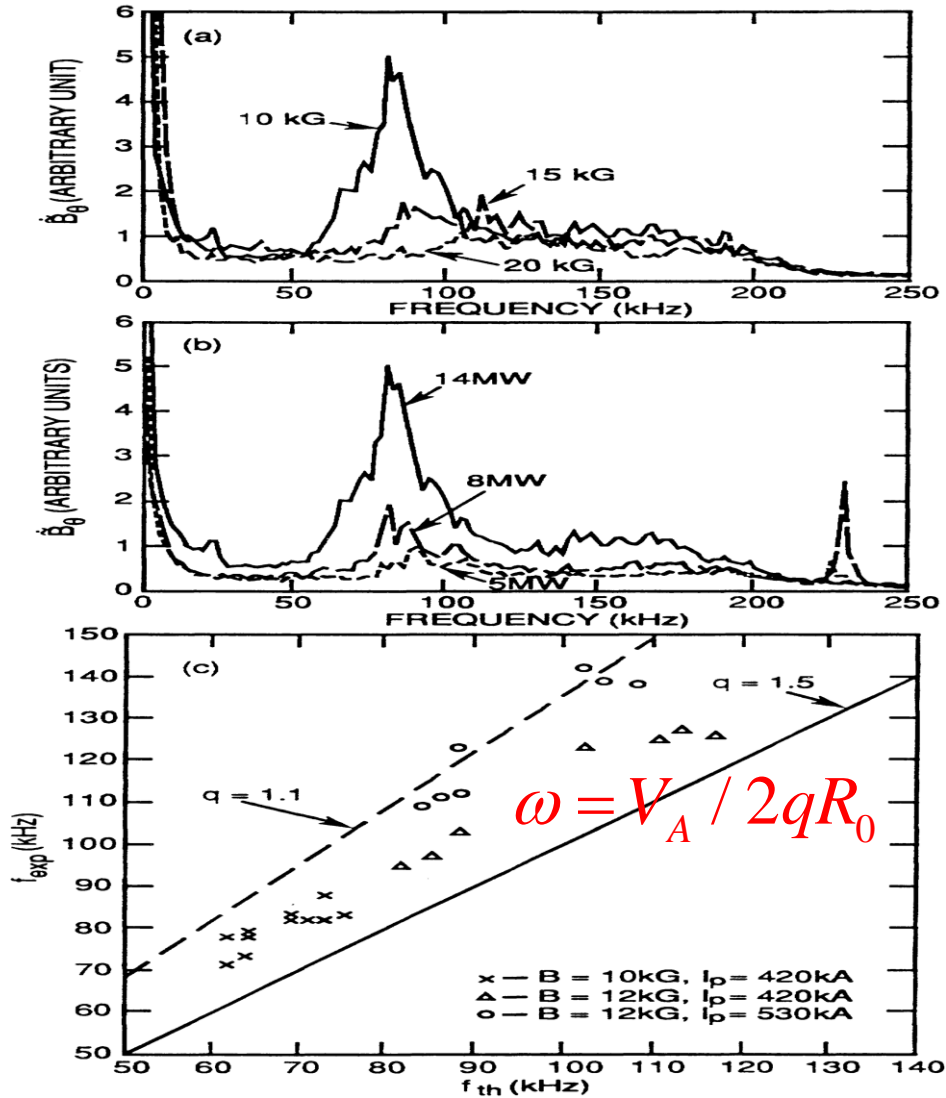
Macor et al PRL 2009

Energetic ion and electron can interact with kink mode to form fishbone.

TAE: Toroidal Alfvén Eigenmode

Wong et al PRL 1991

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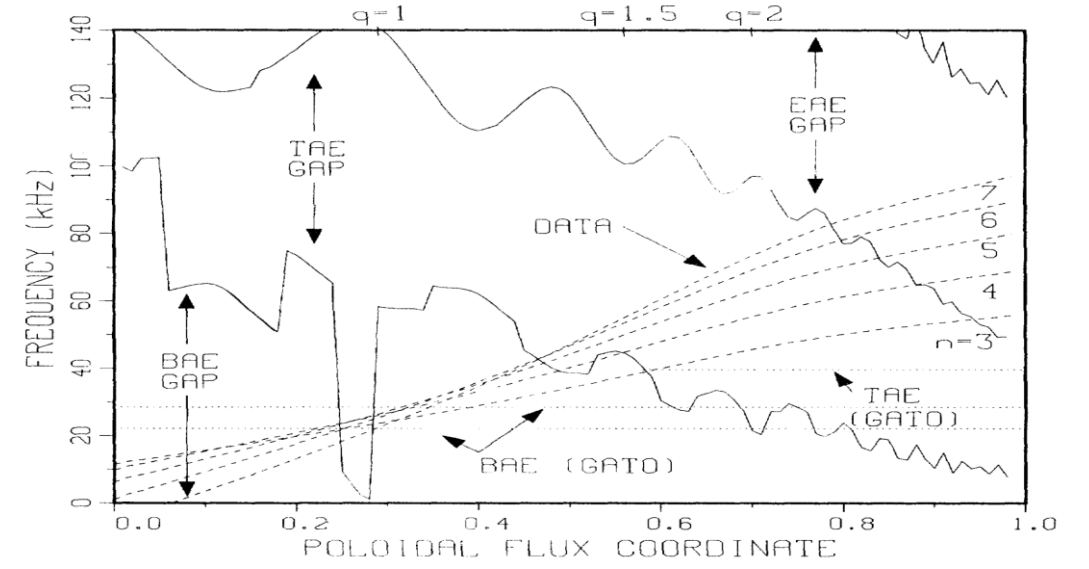
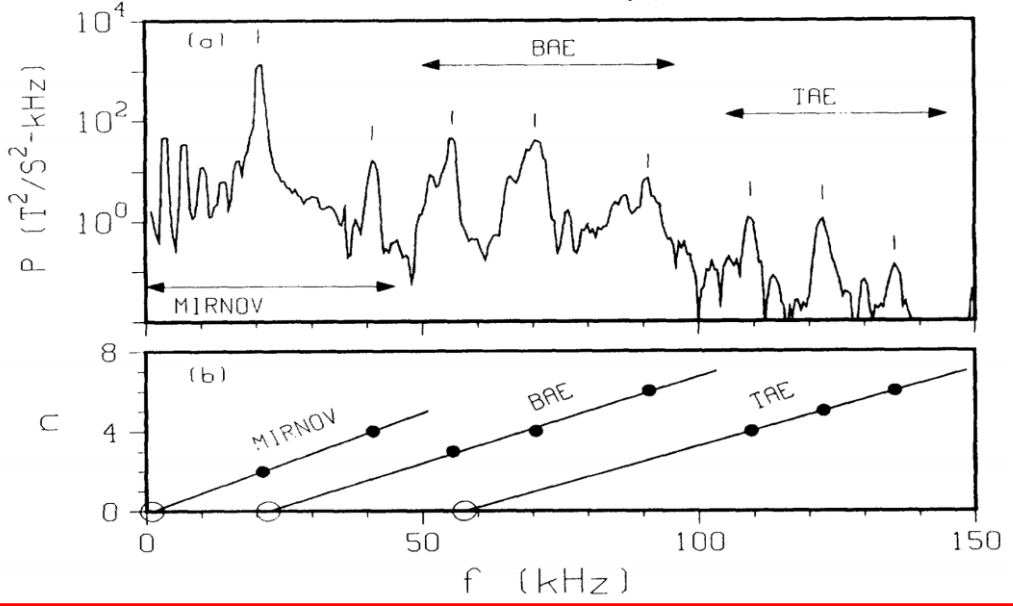
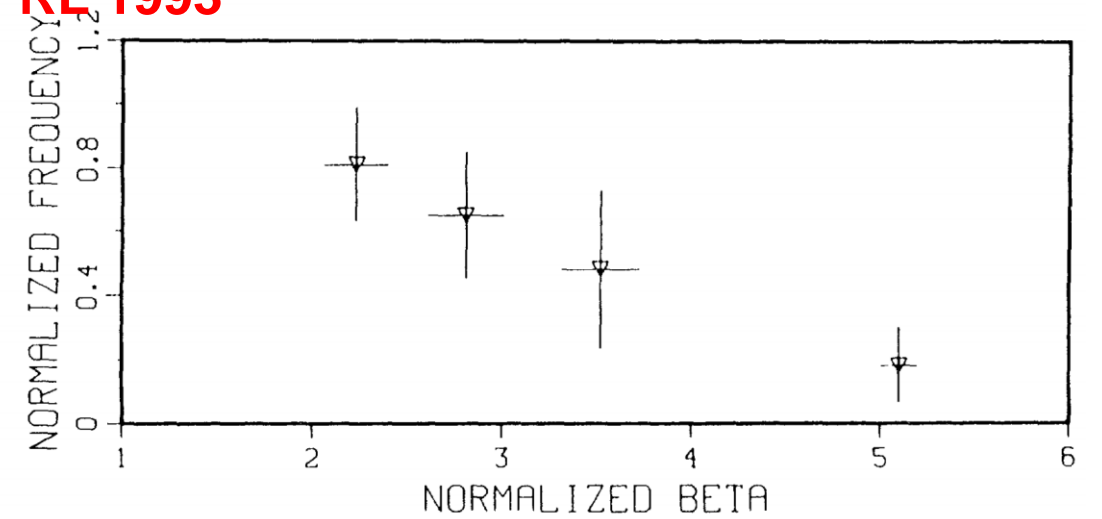
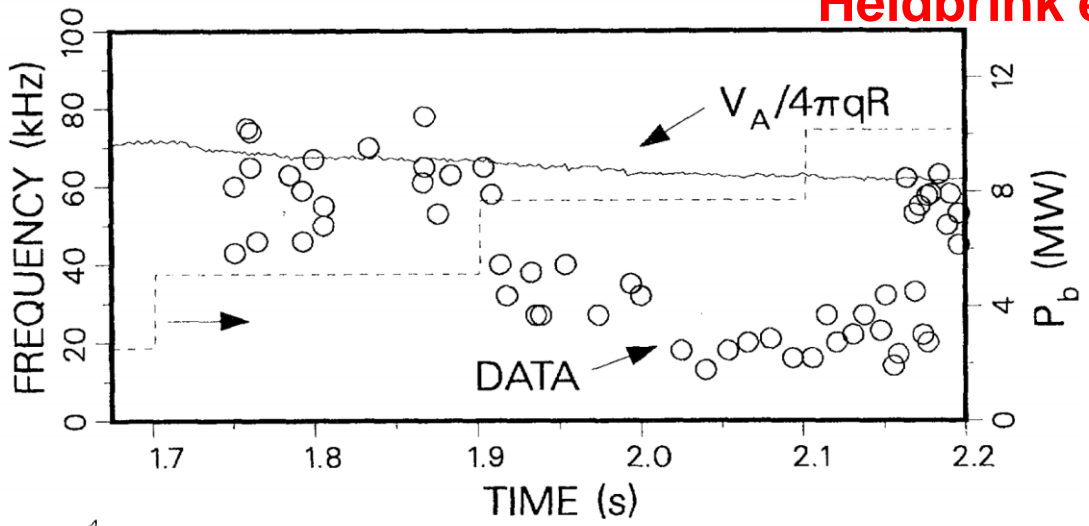


TAE was observed by Mirnov coil and beam emission spectrometry with NBI on TFTR.

BAE: Beta induced Alfvén Eigenmode

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Heidbrink et al PRL 1993



BAE frequency is roughly half of TAE and decreases with β_N based on DIII-D experiments.

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Kinetic-MHD model in NIMROD

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$$\frac{\partial n_\alpha}{\partial t} + \nabla \cdot (n \mathbf{V})_\alpha = \nabla \cdot \mathbf{D} \nabla n_\alpha$$

$$\rho \left(\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} \right) = \mathbf{J} \times \mathbf{B} - \nabla p$$

$$+ \nabla \cdot \rho \nu \nabla \mathbf{V} - \nabla \cdot \mathbf{\Pi} - \nabla p_h$$

$$\frac{n_\alpha}{\Gamma - 1} \left(\frac{\partial \mathbf{T}_\alpha}{\partial t} + \mathbf{V}_\alpha \cdot \nabla \mathbf{T}_\alpha \right) = -p_\alpha \nabla \cdot \mathbf{V}_\alpha$$

$$- \nabla \cdot \mathbf{q}_\alpha + \mathbf{Q}_\alpha - \mathbf{\Pi}_\alpha : \nabla \mathbf{V}$$

- Resistive MHD
- Hall and 2-fluid
- Braginski and beyond closures
- Energetic particles

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times \mathbf{E} + \kappa_{divb} \nabla (\nabla \cdot \mathbf{B})$$

$$\mathbf{J} = \nabla \times \mathbf{B}$$

$$\mathbf{E} = -\mathbf{V} \times \mathbf{B} + \eta \mathbf{J}$$

$$+ \frac{m_e}{n_e e^2} \left[\frac{m_e}{\mu_0} (\mathbf{J} \times \mathbf{B} - \nabla p_e) \right]$$

$$+ \frac{\partial \mathbf{J}}{\partial t} + \nabla (\mathbf{J} \mathbf{V} + \mathbf{V} \mathbf{J})]$$

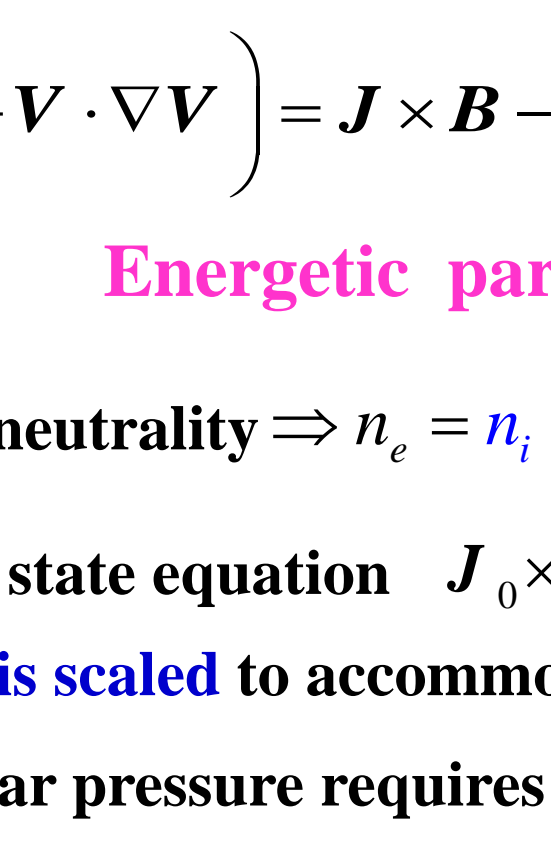
Sovinec et al JCP 2004; Kim et al CPC 2004

Kinetic-MHD model in NIMROD

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- Momentum equation modified by **hot particle pressure tensor**:

$$\rho \left(\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} \right) = \mathbf{J} \times \mathbf{B} - \nabla P_b - \nabla \cdot \mathbf{P}_h + \nabla \cdot \rho \mathbf{v} \nabla \mathbf{V} - \nabla \cdot \mathbf{\Pi}$$



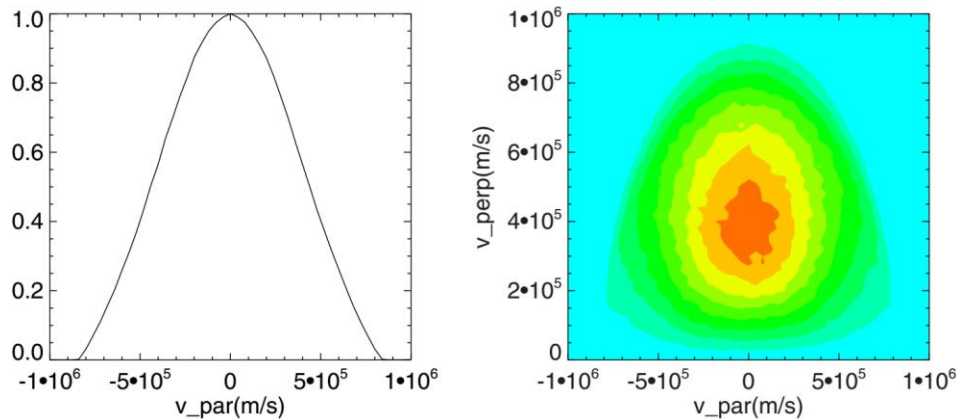
- Quasi-neutrality $\Rightarrow n_e = n_i + n_h$
- Steady state equation $\mathbf{J}_0 \times \mathbf{B}_0 = \nabla p_0 = \nabla p_{b0} + \nabla p_{h0}$
 - ∇p_{b0} is scaled to accommodate hot particles
 - Scalar pressure requires **isotropic** velocity distribution
- Alternative \mathbf{J}_h current coupling possible

Kim et al POP 2008

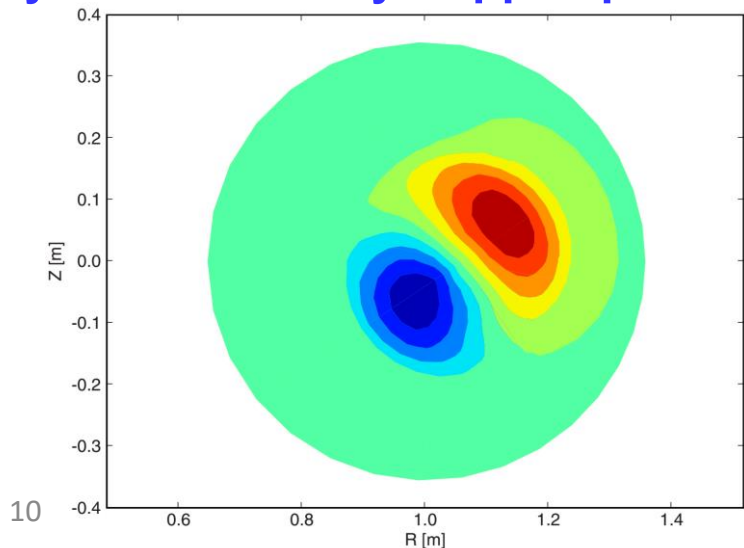
Kink mode driven by slowing down EP

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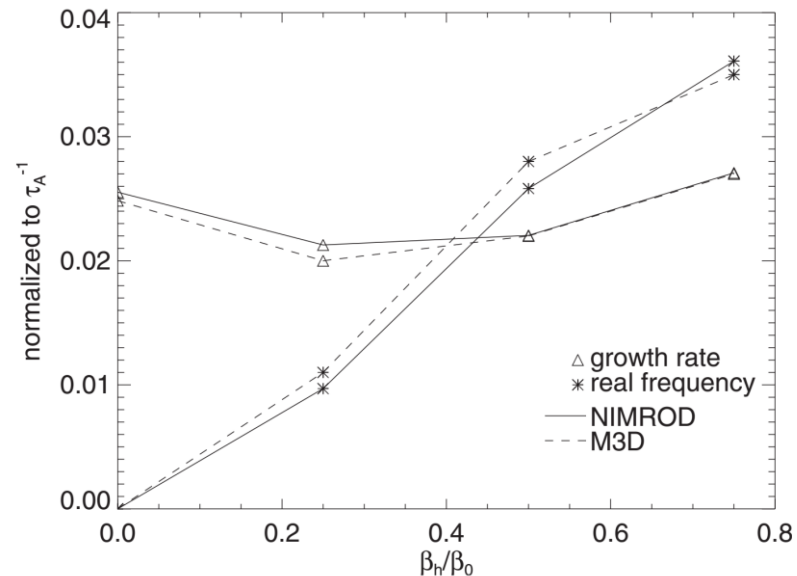
Initial slowing down EP distribution



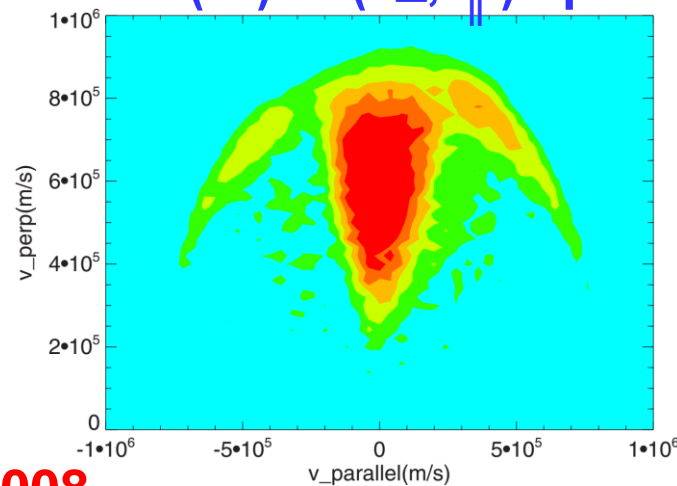
$n=1$ anisotropic pressure contour
mainly contributed by trapped particles



Benchmarked with M3D/M3D-k



ABS(δf) in $(v_{\perp}, v_{\parallel})$ space



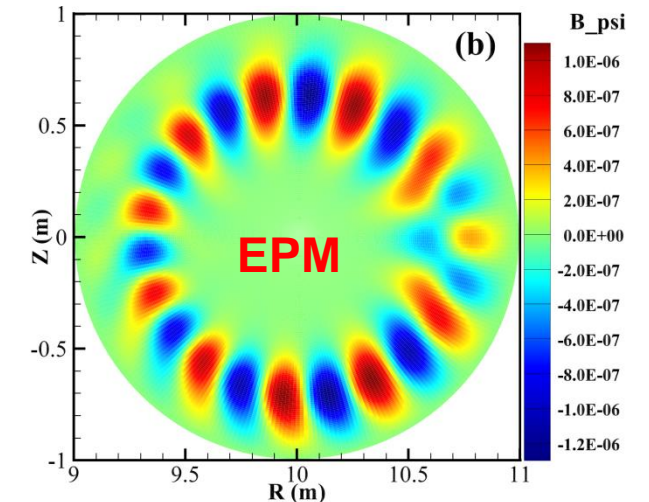
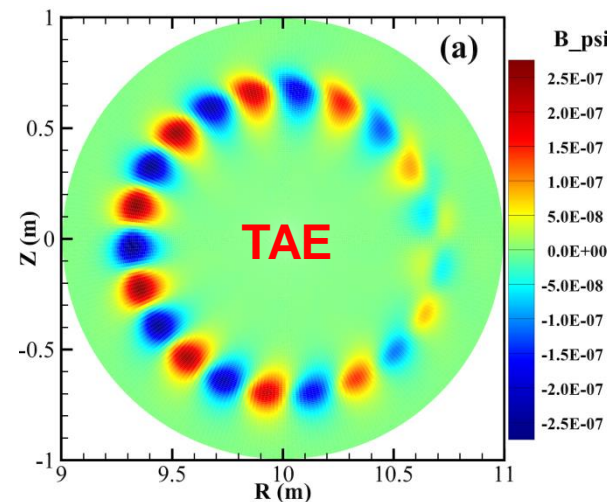
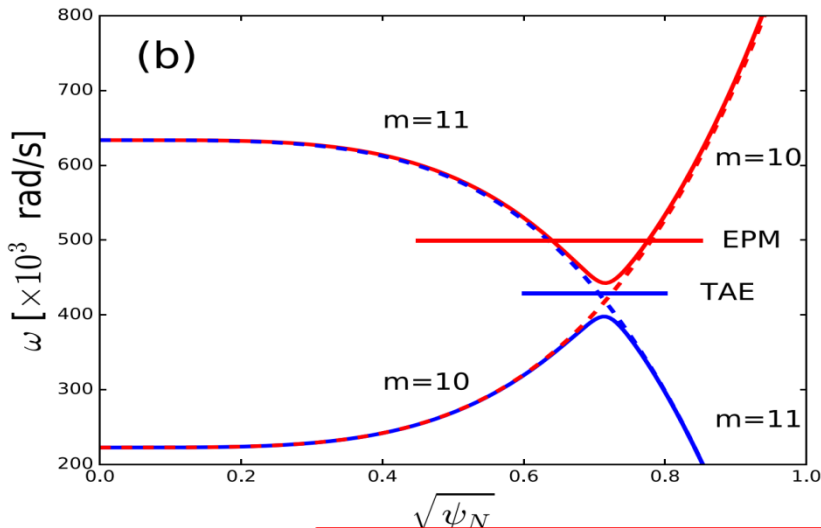
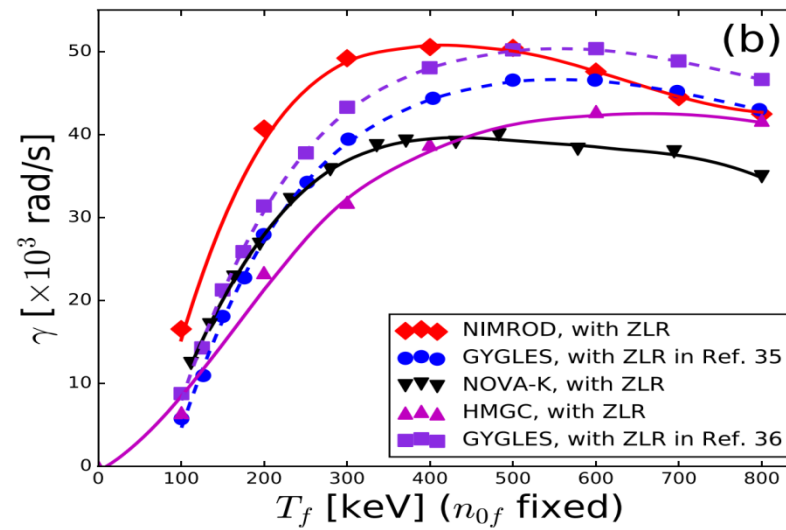
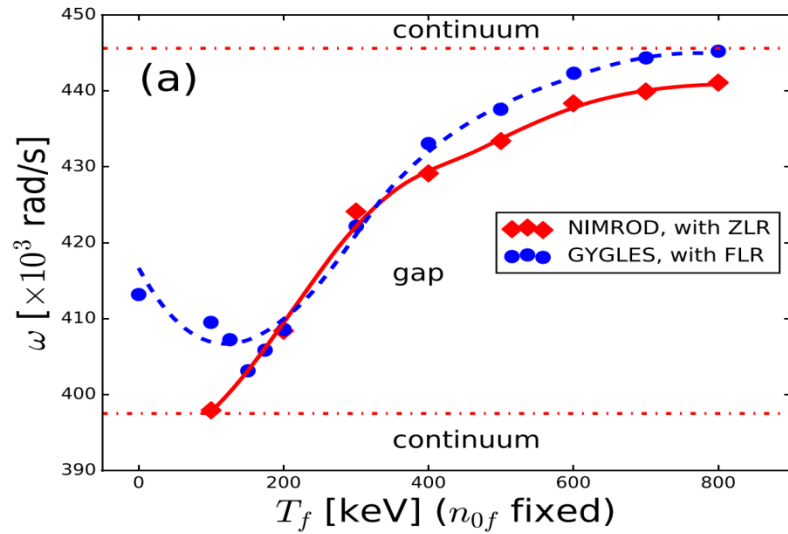
Kim et al POP 2008

TAE driven by Maxwellian EP using NIMROD

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Hou et al POP 2018

Well benchmarked with NOVA-K, HMGC and GYGLES.

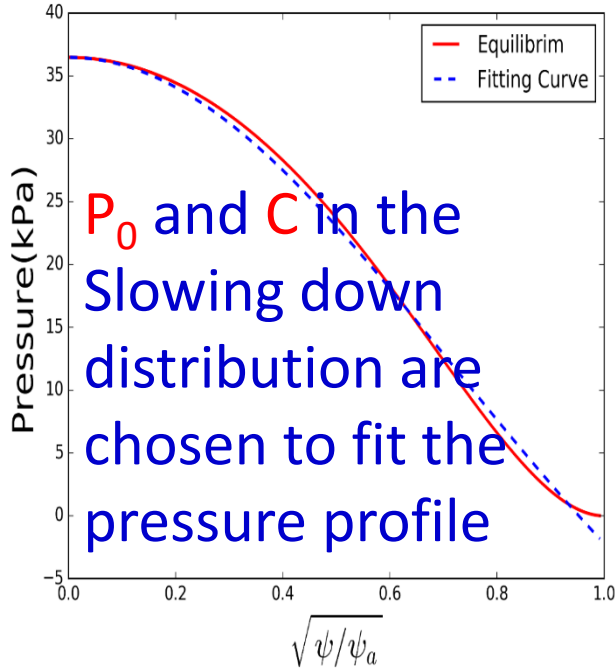
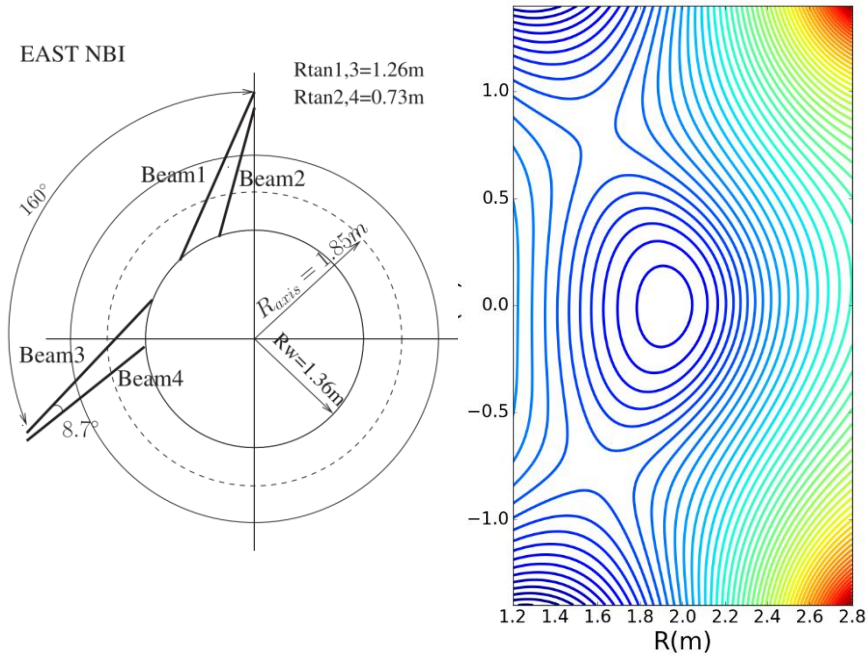


Transition from TAE to EPM occurs as EP β_f increases.

EAST shot #48916 & energetic particle model in NIMROD

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Hou et al POP 2019

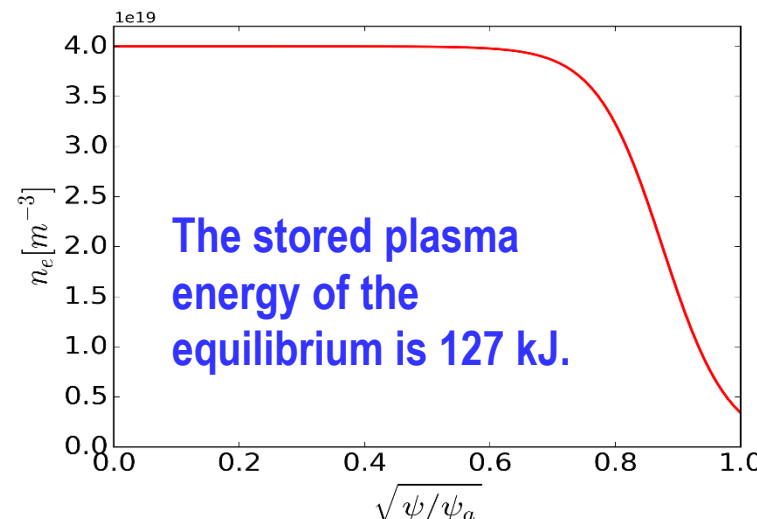
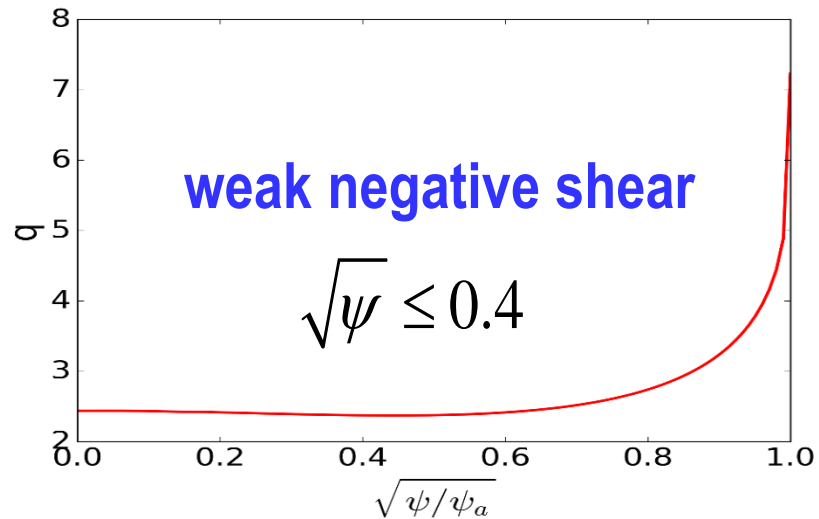


- Energetic particles D^+ is set to be slowing down distribution

$$f_h = \frac{P_0 \exp(g \rho_{||} - \psi_p / c\psi_0)}{\varepsilon^{3/2} + \varepsilon_c^{3/2}}$$

$$g = RB_\phi \quad \rho_{||} = \frac{mv_{||}}{qB}$$

- Critical energy $\varepsilon_c = 37.3keV$
- Beta fraction $\beta_h / \beta = 0.2835$
- EP stored energy: 36kJ

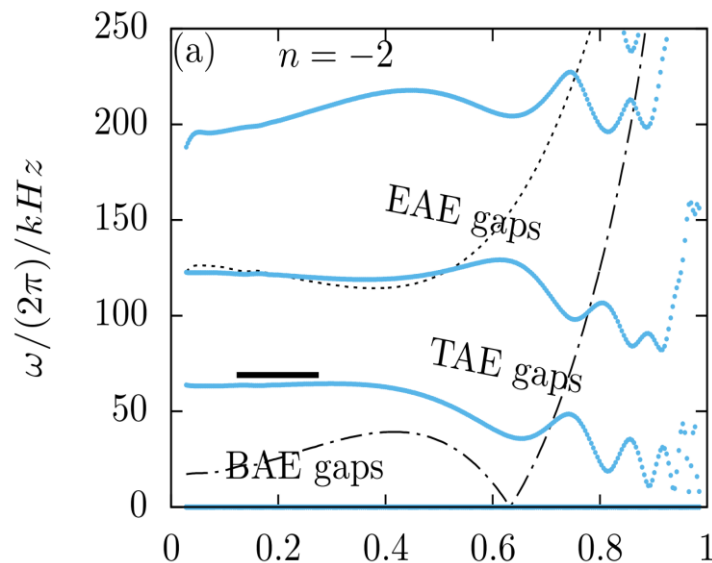


Continuum & mode structure - n=2

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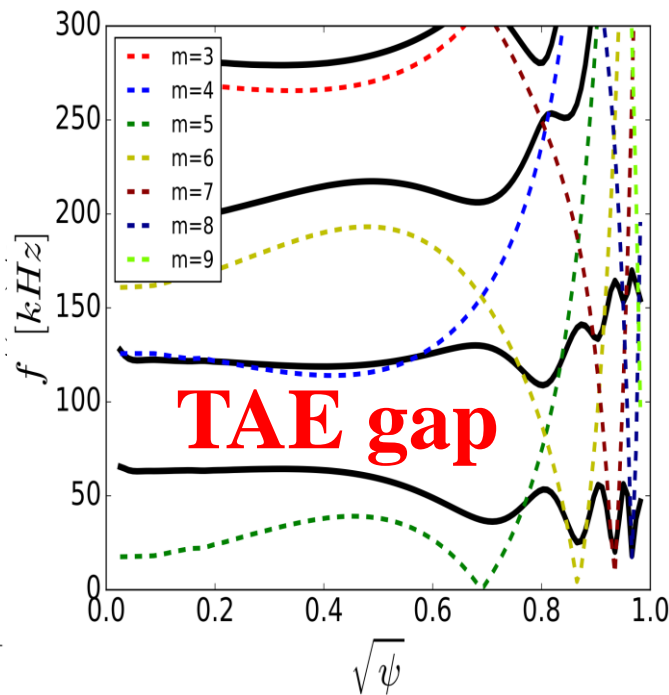
GTAW

cylindrical m=4
 cylindrical m=5 - - - -
 RSAE $f=69\text{kHz}$ ———



Hu et al POP 2016 $\sqrt{\Psi_t}$

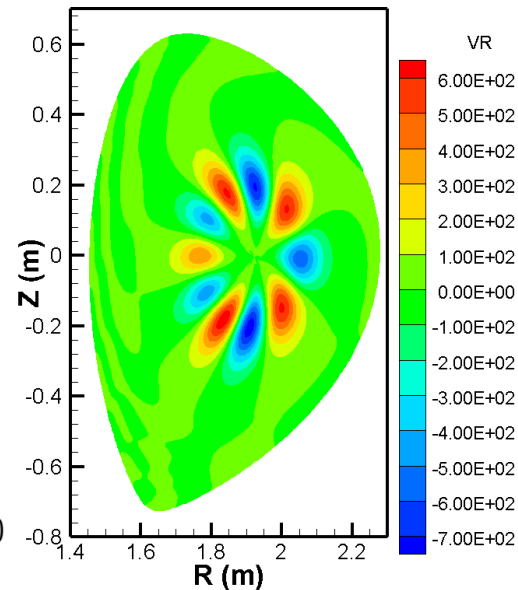
AWEAC



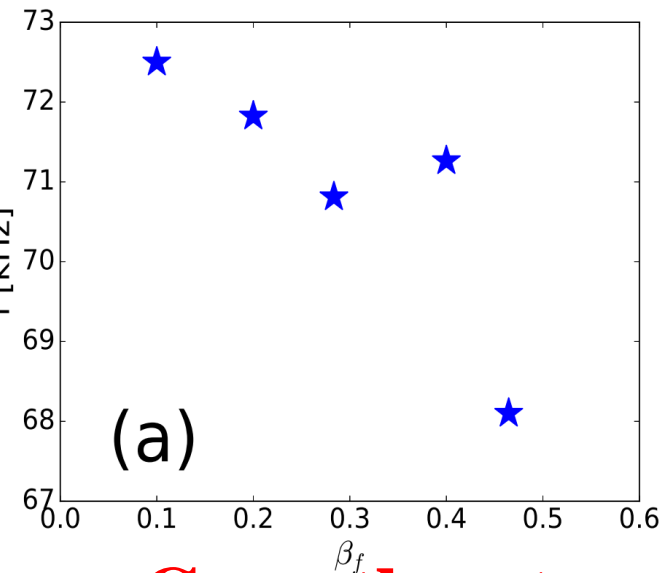
NIMROD

$\gamma = 2.395 \times 10^4 \text{ s}^{-1}$,

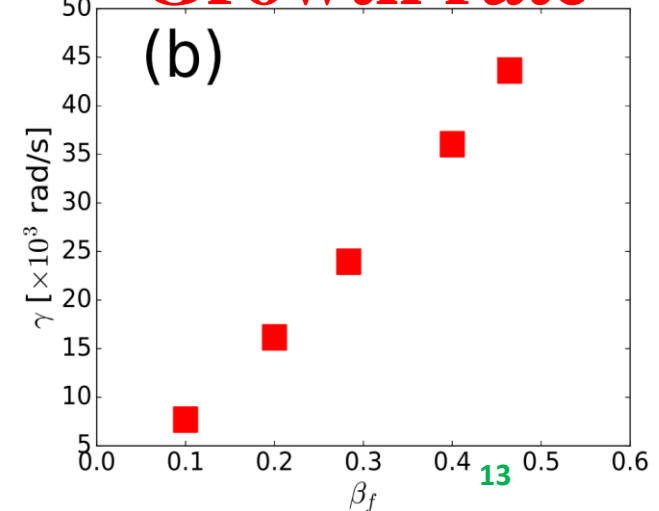
$f = 70.8\text{kHz}$



Frequency



Growth rate



TAE gap is induced by interaction of m=4 & 5 in radial range 0.2-0.4. RSAE with frequency around 70kHz is found from both eigen analysis and NIMROD simulation.

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AWEAC: Alfvén Wave Eigen-Analysis Code

□ Main functions:

- General tokamak geometry (asymmetric D shape); (achieved)
- Read equilibrium data (EFIT, TOQ, NIMEQ, etc.), output the characteristic parameters and visualizations; (achieved)
- Output the Alfvén continuum (BAEs, TAEs, EAEs,...) with different mode number; (achieved)
- Achieve eigenmode growth rate and different damping rates, similar to NOVA/NOVA-k; (under development)
- Density effect due to fueling; flow effect; (under preparation)
- Platform for eigen-analysis in tokamaks. (future goal)

Ideal MHD eigenmode equations

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- Equilibrium eqs.

$$\mathbf{J} \times \mathbf{B} = \nabla P$$

$$\nabla \times \mathbf{B} = \mathbf{J}$$

$$\nabla \cdot \mathbf{B} = \mathbf{0}$$

- Linearized ideal MHD eqs.

$$p_1 + \xi \cdot \nabla P + \gamma_s P \nabla \cdot \xi = 0$$

$$\rho \omega^2 \xi = \nabla p_1 + \mathbf{b} \times \nabla \times \mathbf{B} + \mathbf{B} \times \nabla \times \mathbf{b}$$

$$\mathbf{b} = \nabla \times (\xi \times \mathbf{B})$$



- Ideal MHD eigenmode eqs.

$$\nabla \psi \cdot \nabla \begin{pmatrix} P_1 \\ \xi_\psi \end{pmatrix} = C \begin{pmatrix} P_1 \\ \xi_\psi \end{pmatrix} + D \begin{pmatrix} \xi_s \\ \nabla \cdot \xi \end{pmatrix}$$

$$E \begin{pmatrix} \xi_s \\ \nabla \cdot \xi \end{pmatrix} = F \begin{pmatrix} P_1 \\ \xi_\psi \end{pmatrix}$$

- Where

$$P_1 = p_1 + \mathbf{b} \cdot \mathbf{B}$$

$$\xi_s = \xi \cdot (\mathbf{B} \times \nabla \psi) / |\nabla \psi|^2$$

$$\xi_\psi = \xi \cdot \nabla \psi$$

Alfven wave continuum equations

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- Ideal MHD Alfven continuum eq.

$$E \begin{pmatrix} \xi_s \\ \nabla \cdot \xi \end{pmatrix} = \begin{pmatrix} E_{11} & E_{12} \\ E_{12} & E_{22} \end{pmatrix} \begin{pmatrix} \xi_s \\ \nabla \cdot \xi \end{pmatrix} = 0$$



$$E = E_a + \omega^2 E_b$$

$$E_a \begin{pmatrix} \xi_s \\ \nabla \cdot \xi \end{pmatrix} = -\omega^2 E_b \begin{pmatrix} \xi_s \\ \nabla \cdot \xi \end{pmatrix}$$

- Where

$$E_{11}^a = -\mu_0^{-1} \mathbf{B}_0 \cdot \nabla \left(\frac{|\nabla \psi|^2 \mathbf{B}_0 \cdot \nabla}{B_0^2} \right)$$

$$E_{12}^a = -2\gamma P_0 \kappa_s,$$

$$E_{21}^a = 0,$$

$$E_{22}^a = \frac{\gamma P_0}{-\mu_0 \rho_0} \mathbf{B}_0 \cdot \nabla \left(\frac{\mathbf{B}_0 \cdot \nabla}{B_0^2} \right)$$

$$E_{11}^b = -\frac{\rho_0 |\nabla \psi|^2}{B_0^2}$$

$$E_{12}^b = 0,$$

$$E_{21}^b = \kappa_s$$

$$E_{22}^b = \frac{\gamma P_0 + \mu_0^{-1} B_0^2}{B_0^2}$$

$$\kappa = \mathbf{b}_0 \cdot \nabla \mathbf{b}_0$$

$$= (\nabla \times \mathbf{b}_0) \times \mathbf{b}_0$$

$$\kappa_s = 2\kappa \cdot \frac{\mathbf{B}_0 \cdot \nabla \psi}{B_0^2}$$

- Slow sound approximation

$$\frac{E_{22}^a}{E_{22}^b} = \frac{\gamma\beta / 2}{1 + \gamma\beta / 2}$$



$$\beta \ll 1,$$

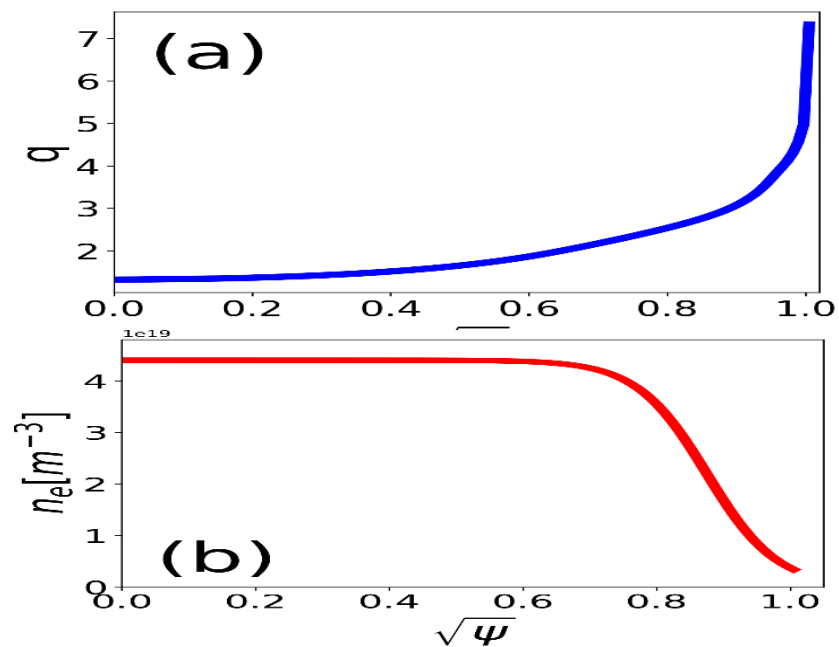
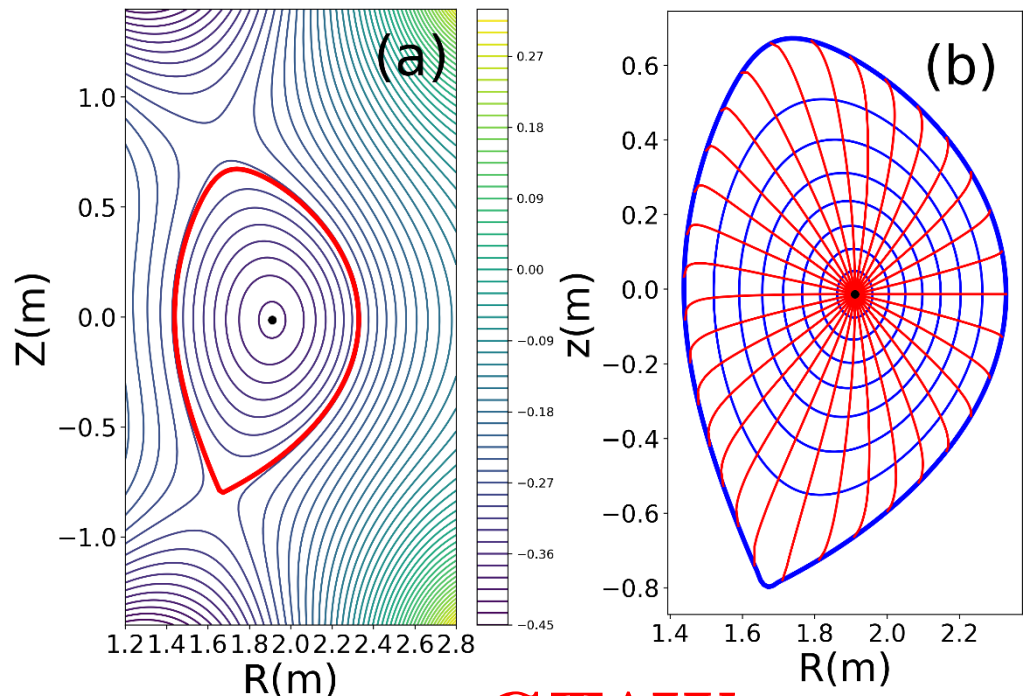
$$E_{22}^a \rightarrow 0$$

Hu et al POP 2014;
Deng et al NF 2012

Benchmark of AWEAC with GTAW

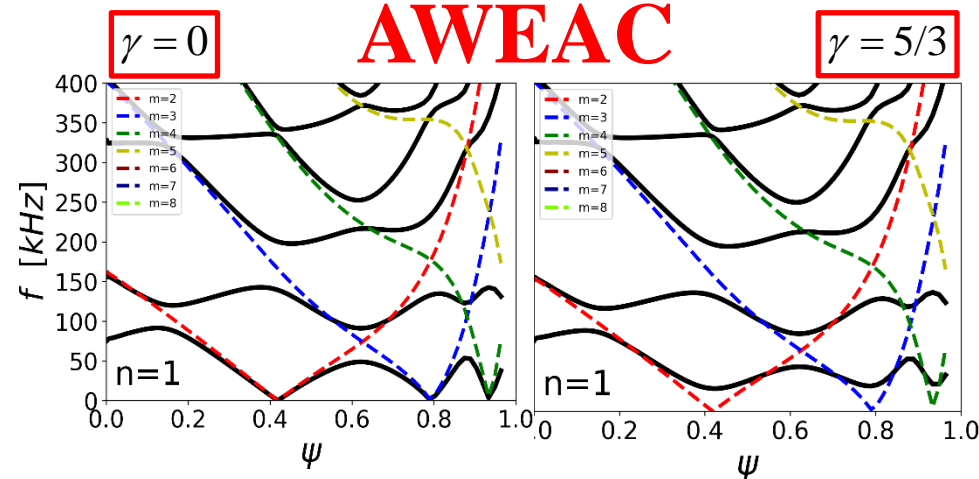
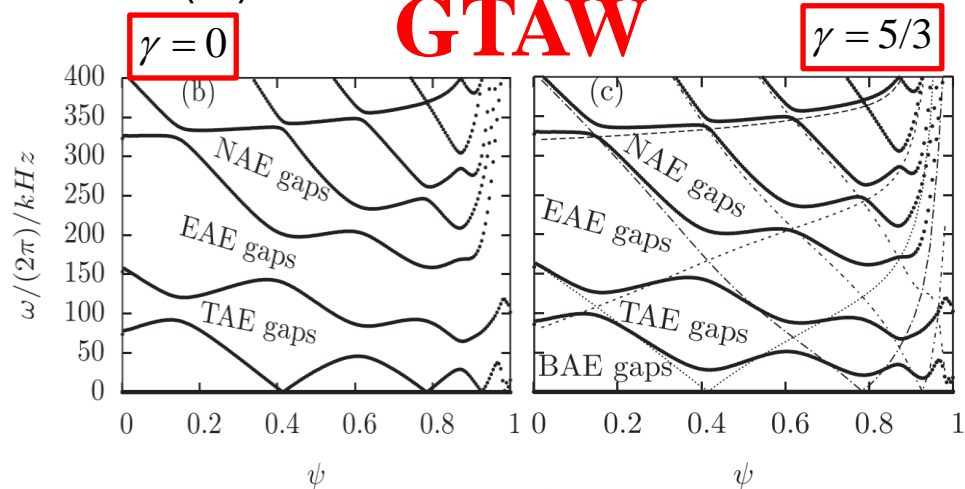
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EAST shot #38300



GTAW

AWEAC



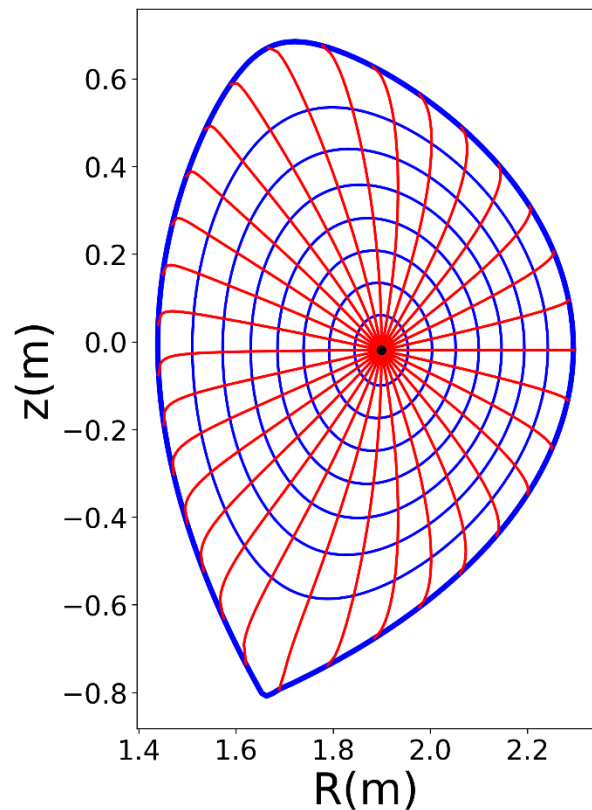
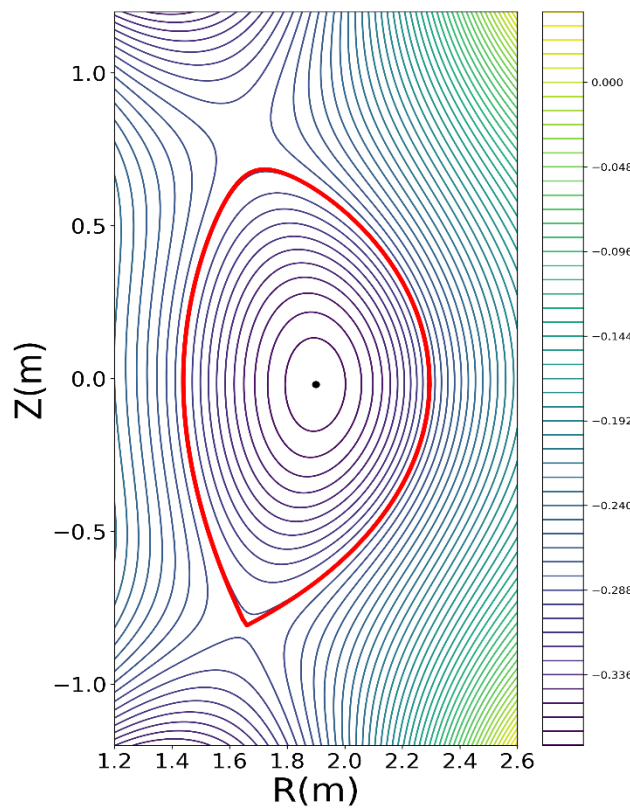
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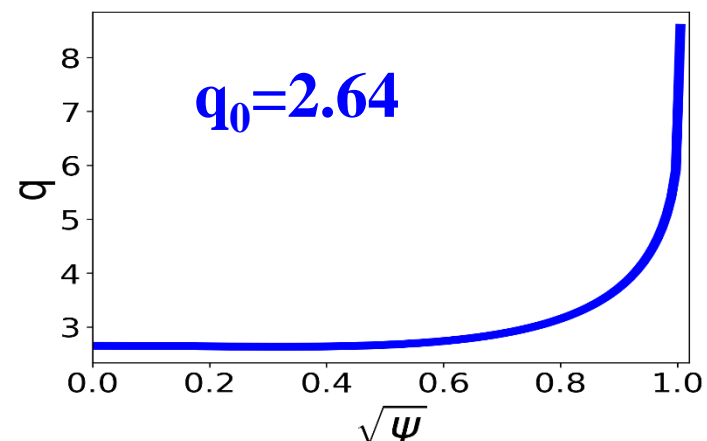
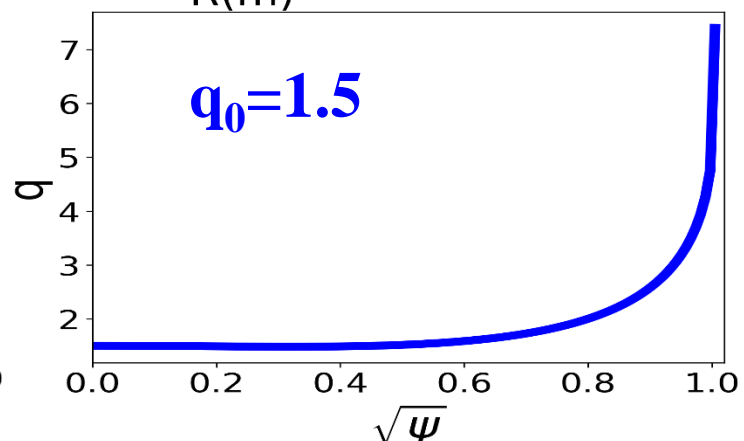
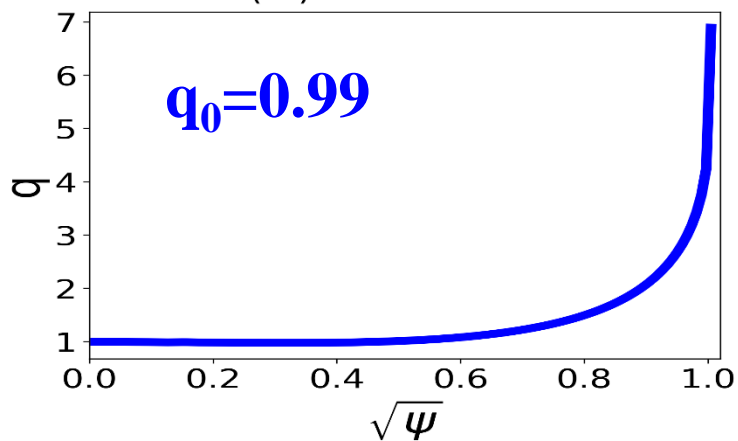
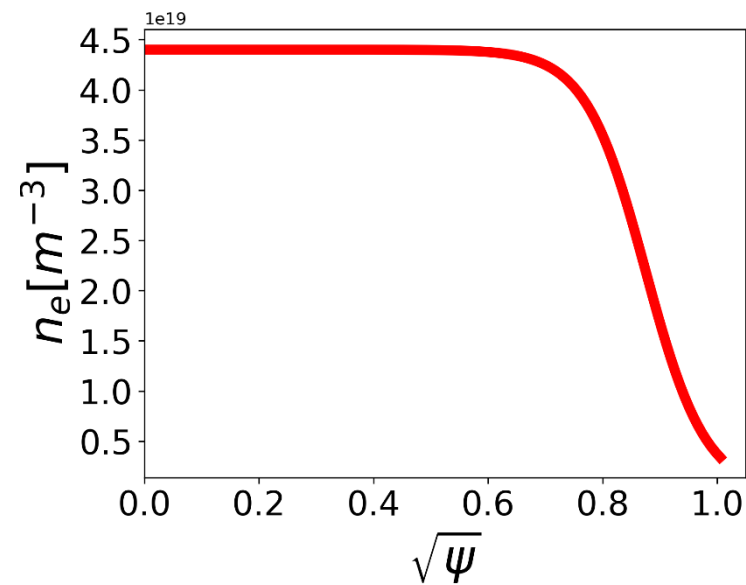
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EP driven MHD modes with varied q_0

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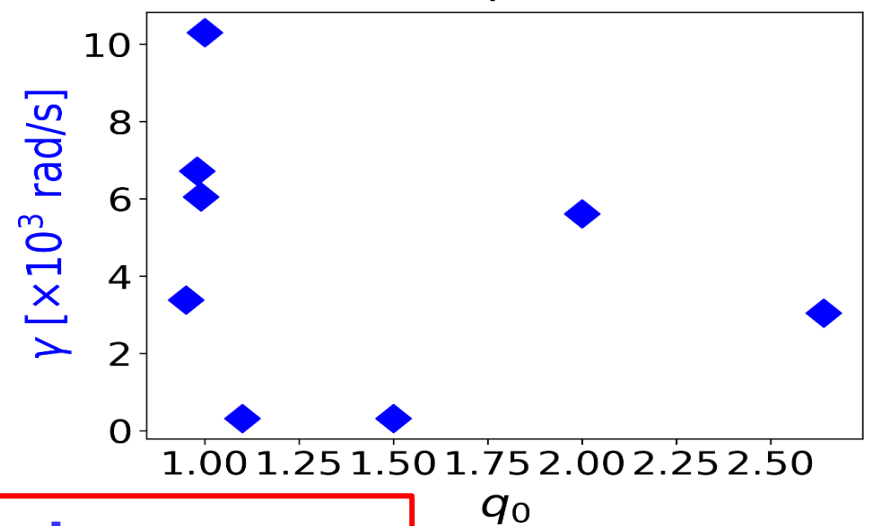
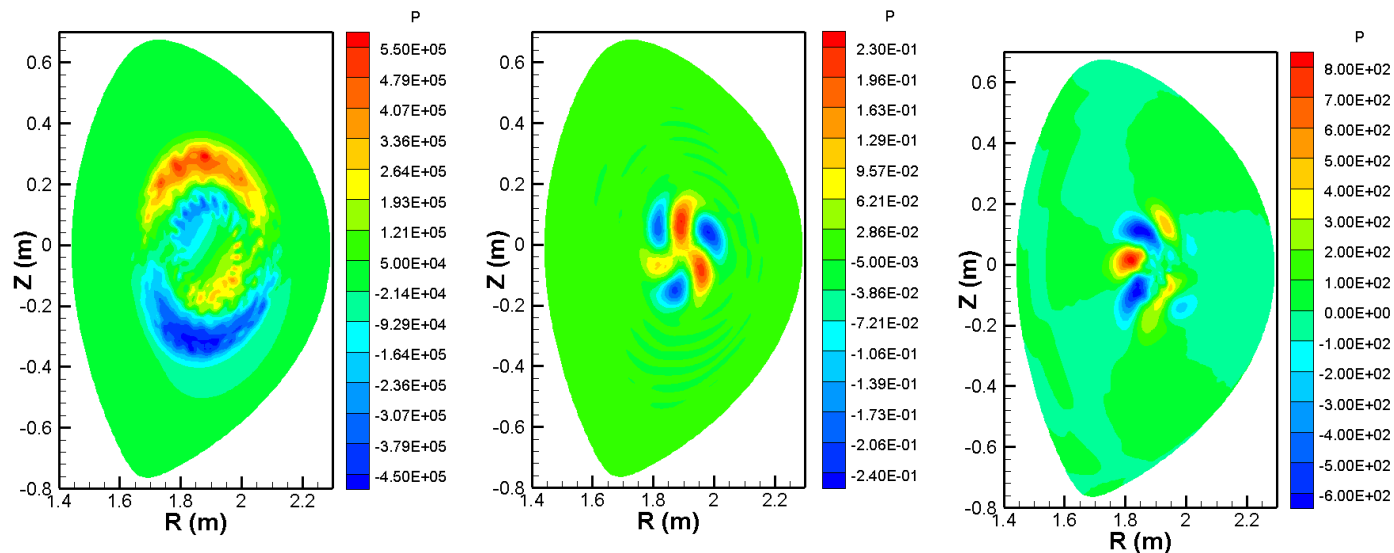
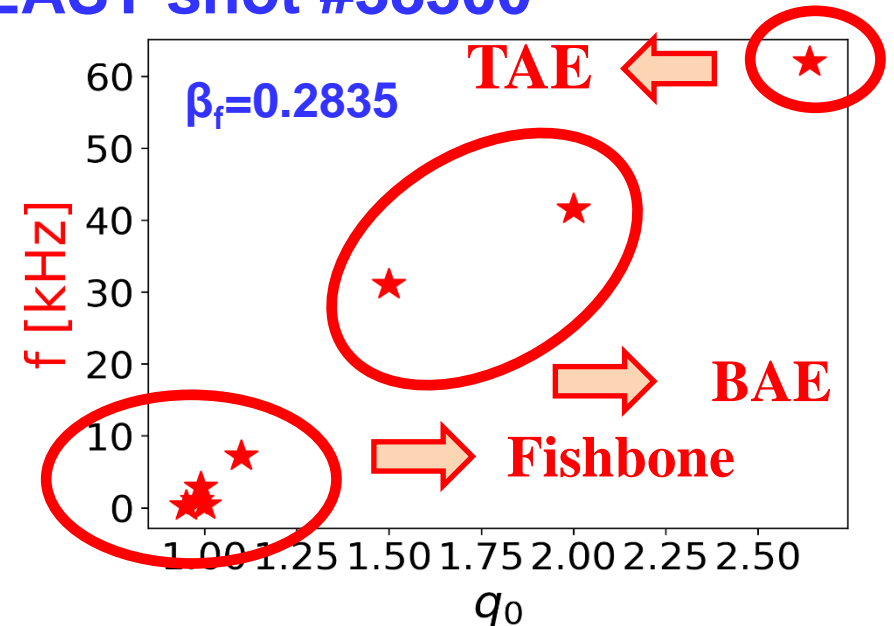
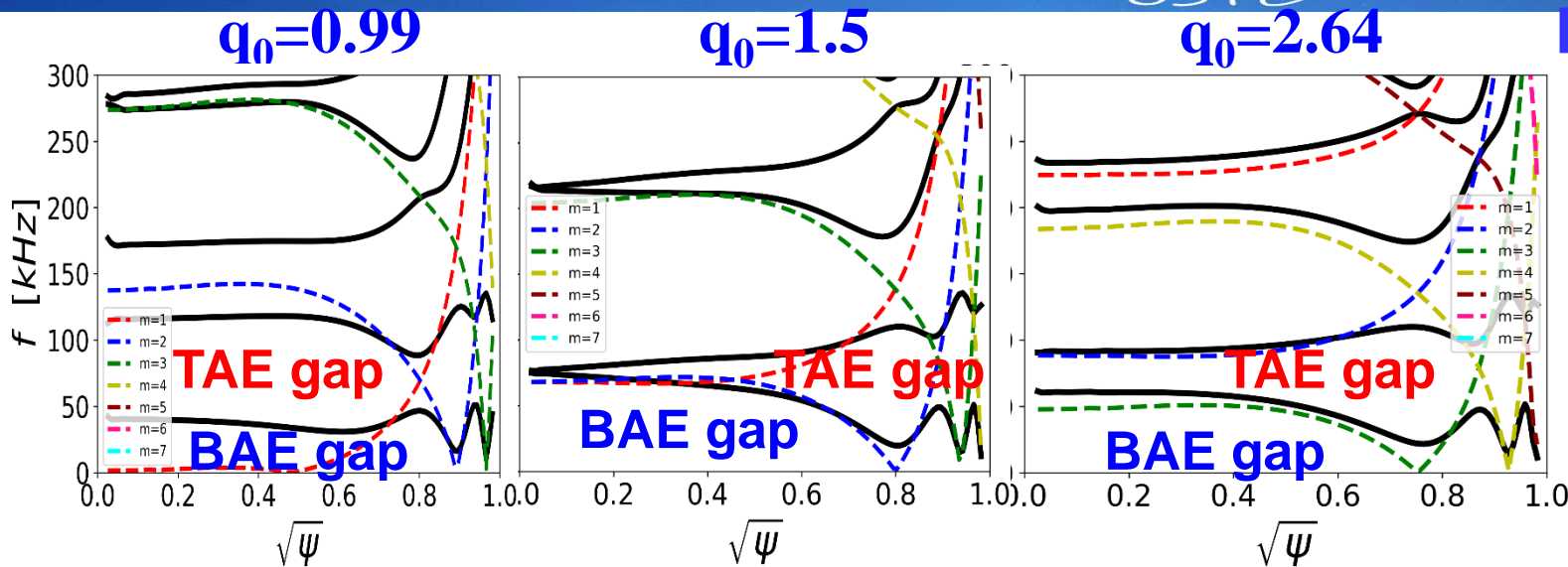
EAST shot #38300



EP driven MHD modes with varied q_0

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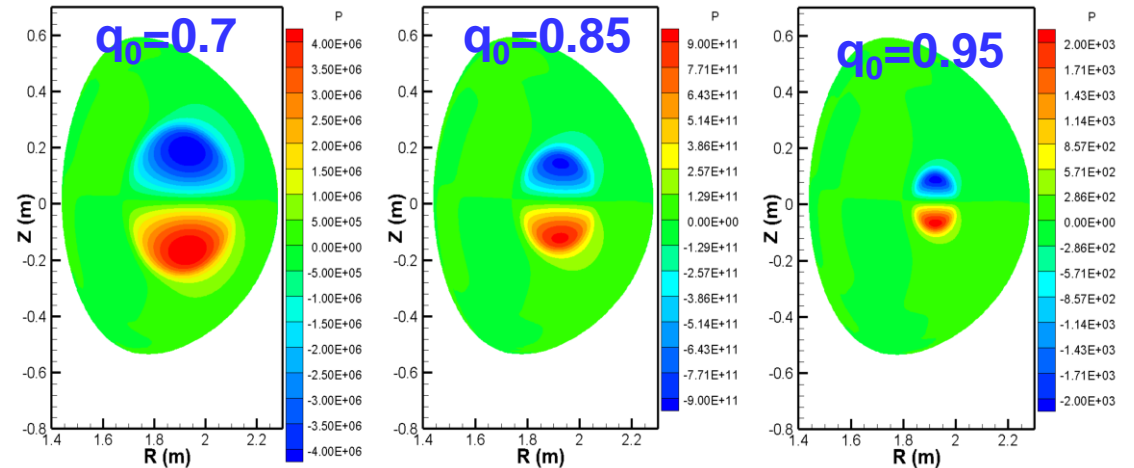
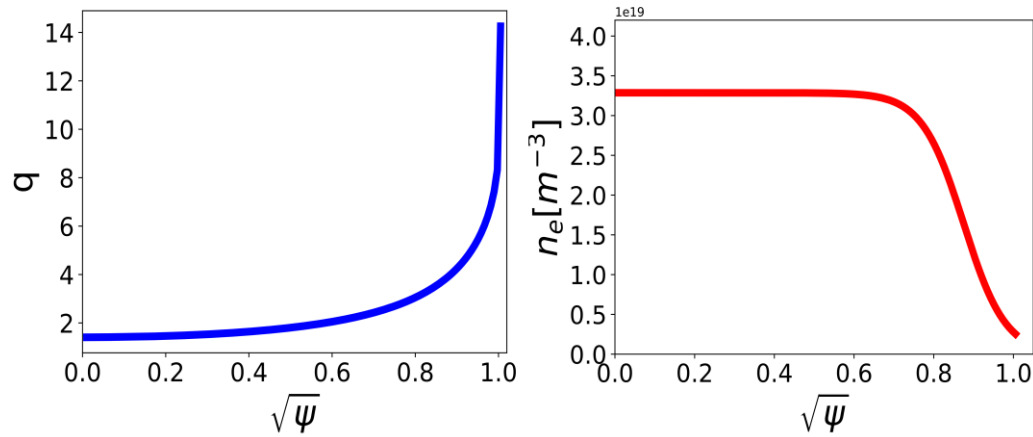
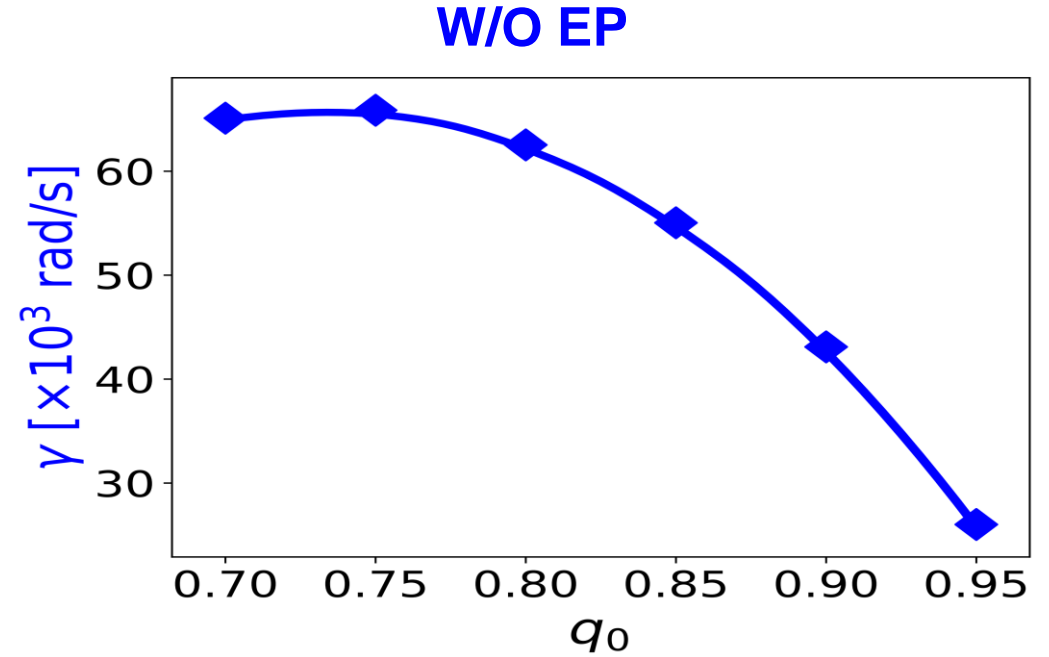
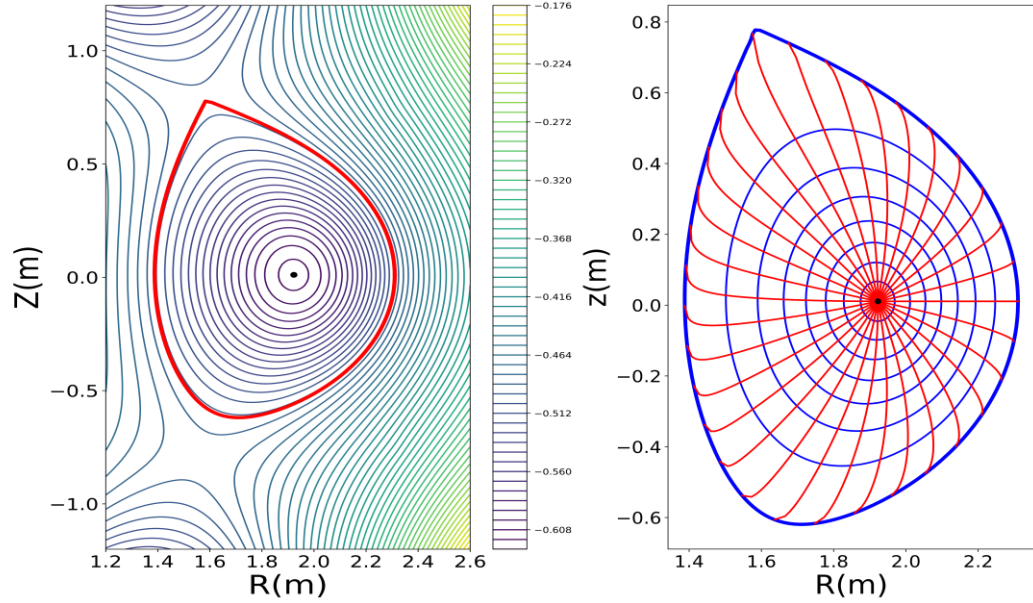
EAST shot #38300



Transition from fishbone to BAE/TAE occurs as q_0 increases.

Kink/fishbone mode study of EAST shot #70187

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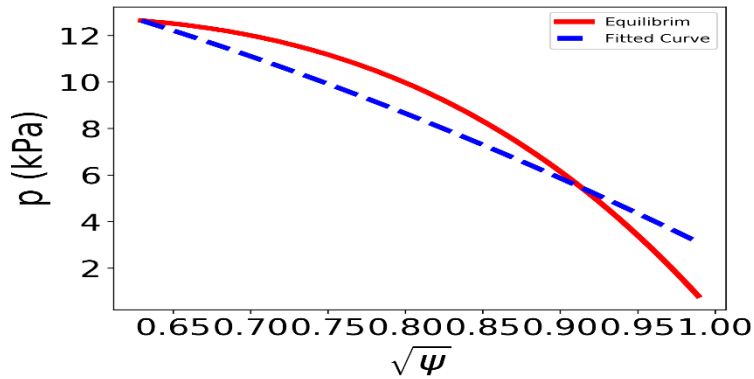


Kink mode increases first and then decreases as q_0 increases.

Kink/fishbone mode study of EAST shot #70187

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With EP

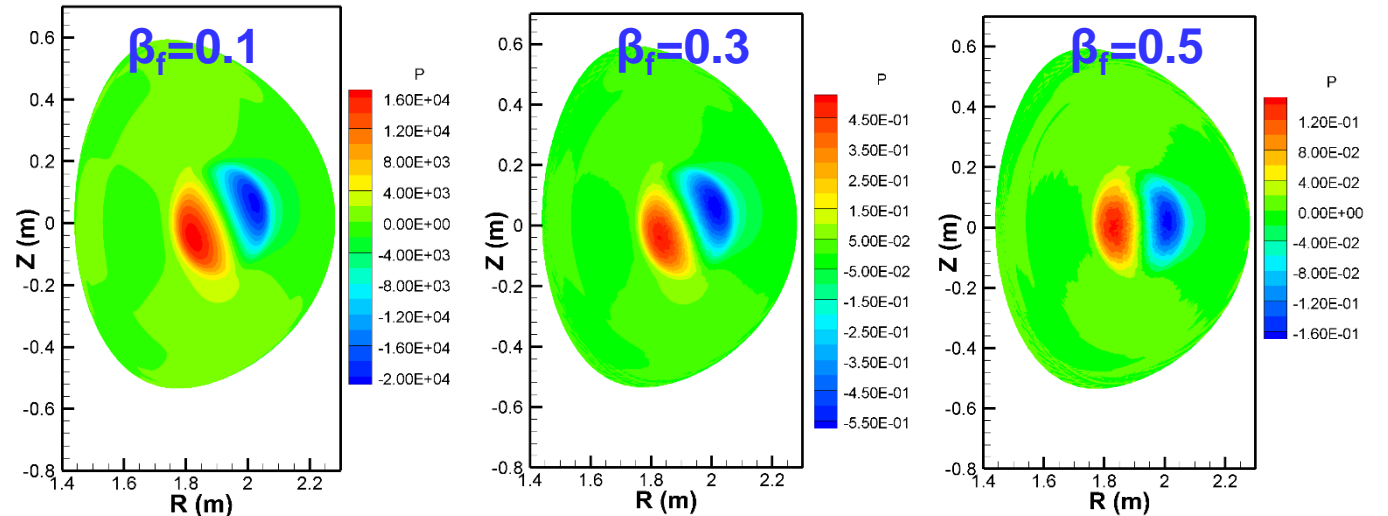
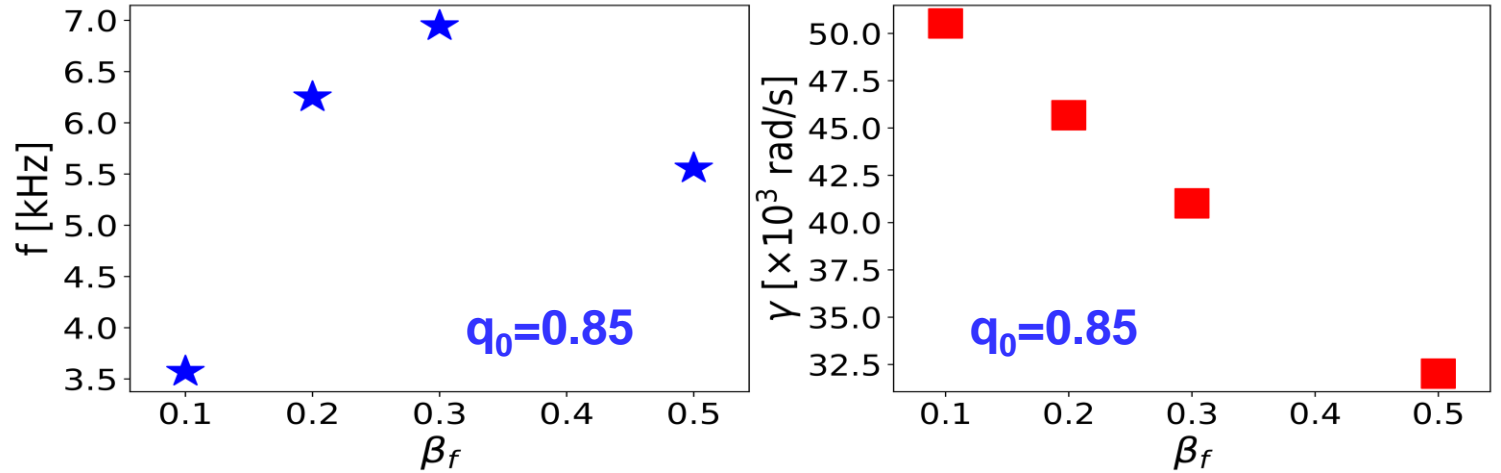


- Energetic particles D+ is set to be slowing down distribution

$$f_h = \frac{P_0 \exp\left(\frac{g \rho_{\parallel} - \psi_p}{c \psi_0}\right)}{\varepsilon^{3/2} + \varepsilon_c^{3/2}}$$

$$g = RB_{\phi} \quad \rho_{\parallel} = \frac{mv_{\parallel}}{qB}$$

- Critical energy $\varepsilon_c = 37.3 \text{ keV}$



Fishbone mode is suppressed as β_f increases.

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Conclusions & future works

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□ Conclusions:

- Kink/fishbone mode & TAE simulation using NIMROD has been benchmarked with both eigenvalue codes and initial value codes for both circular and EAST tokamaks.
- Transition from fishbone to BAE/TAE occurs as q_0 increases for EAST shot #38300;
- BAE is hard to be excited by EP for flat q profile with q_0 less than 1.

□ Future works:

- Simulations of fishbone and AEs on EAST/HL-2A based on experiment parameters, with energetic ions;
- Nonlinear simulations of EP driven AEs or MHD instabilities on EAST/HL-2A/ ITER/CFETR.



中国科学技术大学
University of Science and Technology of China

Thank you!

Suggestions and collaborations are welcome!

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