10th US-PRC Magnetic Fusion Collaboration Workshop (MFCW 2020)



The simulations on the edge instabilities of the RF heating H-mode discharges on EAST

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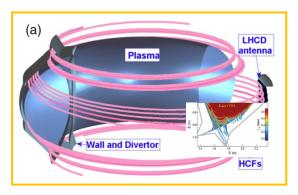
Collaborated with: X.Q. Xu, B. Zhu at LLNL

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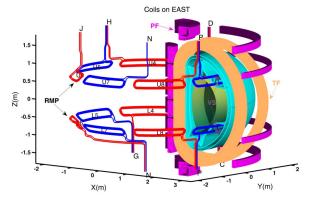


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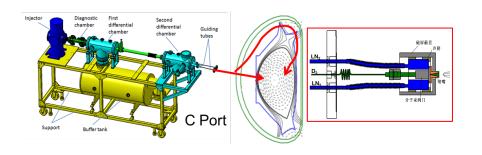
Various effective ELM control methods on EAST



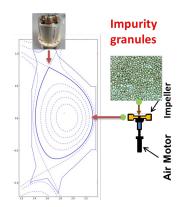
Lower Hybrid wave (LHW) [Liang Y. et al, PRL (2013)]



Resonant magnetic perturbation[Sun Y. et al, PRL (2016)]



Pellet, super-sonic molecular beam injection (SMBI)



Lithium injection [Hu J. et al, PRL (2015)]

- > EAST is the tokamak which can be operated under the ITER-like long-pulse, low rotation and metal wall conditions.
- > Various ways of ELM control have been developed on EAST successfully.



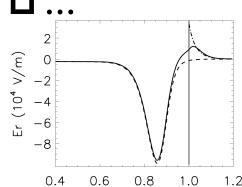
The effects of RF waves on edge instabilities

>Indirect effects:

- ☐ Heating: profiles
- ☐ Current driving: magnetic shear
- □ Ponderomotive force: rotation
- **□**

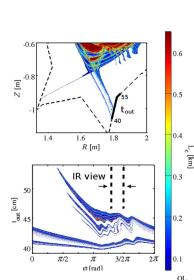
> Direct effects:

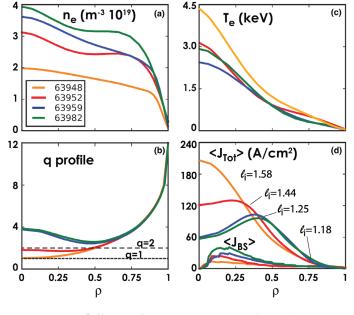
- ☐ Current driving: topology
- ☐ RF sheath: SOL rotation
- ☐ Source terms



Normalized Ψ

B. Gui et al., NF 58 026027 (2018)





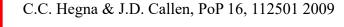
A. Garrofallo et al., NF 57 076037 (2017)

Y.F. Liang et al., PRL 110, 235002 (2013).

$$m_{s}n_{s}\left(\frac{\partial \mathbf{V}_{s}}{\partial t} + \mathbf{V}_{s} \cdot \nabla \mathbf{V}_{s}\right) = n_{s}q_{s}(\mathbf{E} + \mathbf{V}_{s} \times \mathbf{B}) - \nabla p_{s}$$
$$-\nabla \cdot \boldsymbol{\pi}_{s} + \mathbf{R}_{s} + \mathbf{F}_{s0}^{\mathrm{rf}},$$

$$\frac{3}{2}n_s \left(\frac{\partial T_s}{\partial t} + \mathbf{V}_s \cdot \nabla n_s \right) + n_s T_s \nabla \cdot \mathbf{V}_s$$

$$= -\nabla \cdot \mathbf{q}_s - \pi_s : \nabla \mathbf{V}_s + Q_s + S_{s0}^{\mathrm{rf}}.$$

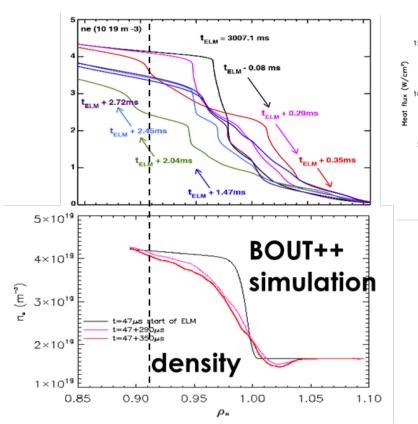




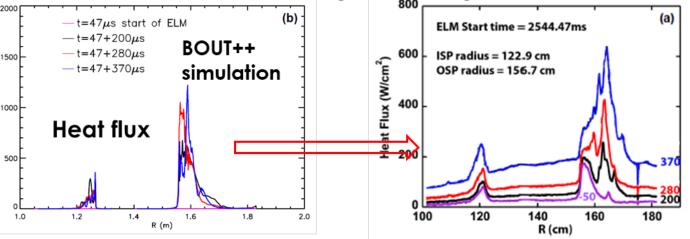


BOUT++ framework has been used to reproduce the profile change & transient fluxes induced by edge instabilities

The 6-field 2-fluid model is widely used on the physics understanding of the edge physics for tokamaks

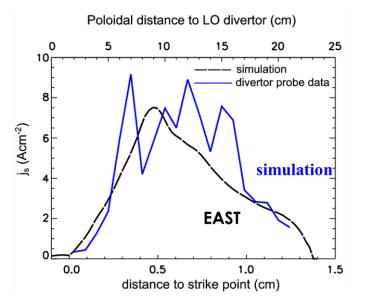


Density evolution during ELM on DIII-D



Transient heat flux during ELM on DIII-D

T.Y. Xia, et al. Nucl. Fusion, 55 (2015) 113030



Transient particle flux by ELM on EAST

Y.B. Wu, T.Y. Xia et al., PCF 60 (2018) 05500

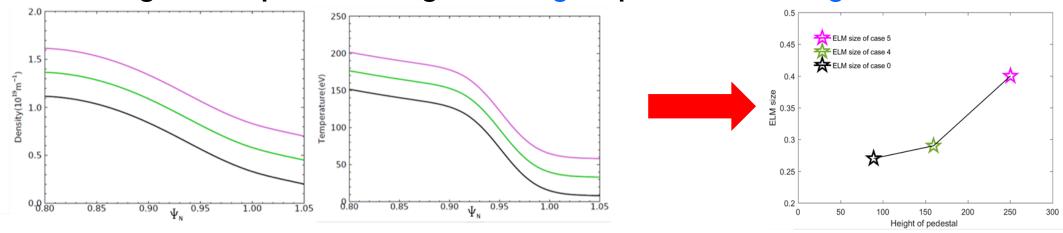


- >Indirect effects by RF:
 - Pedestal scan for ELM
 - Edge turbulence and the interactions with ELM
- ➤ Direct effects by RF:
 - ☐ HCF effects on the turbulence & heat flux
 - ☐ RF sheath on ELM

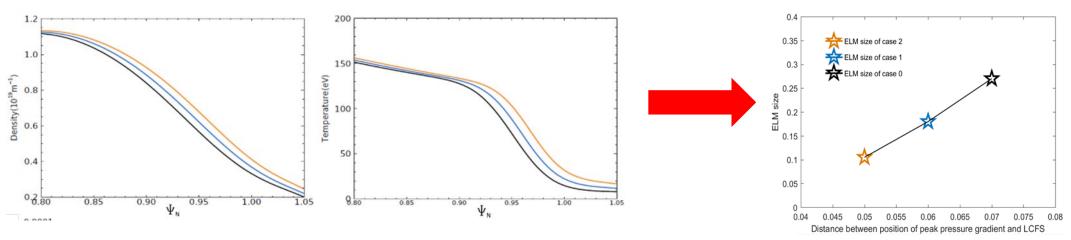


The effects of RF heating on ELMs: pedestal scans

> The change of the pedestal height: The higher pedestal, the larger ELM size



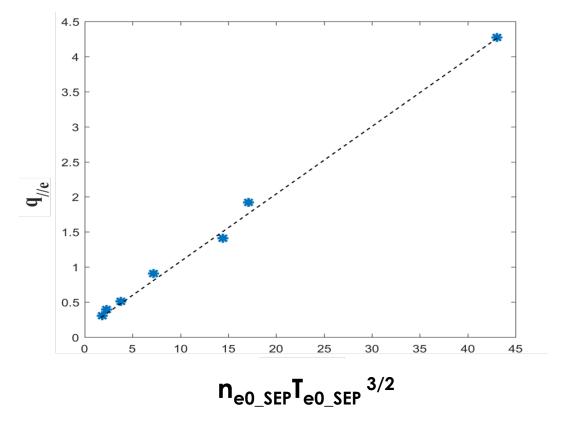
> The change of the pedestal position: The closer of the pedestal to separatrix, the **smaller ELM size**







The peak divertor heat flux is proportional to the upstream density & temperature



Effective thermal conduction in 6-field 2-fluid module of BOUT++

$$\kappa_{\text{eff},j} = \frac{\kappa_{\parallel j} \kappa_{\text{fs},j}}{\kappa_{\parallel j} + \kappa_{\text{fs},j}}$$

$$\kappa_{\mathrm{fs},j} = n_j v_{\mathrm{th},j} q R_0$$

$$q_{\parallel e} \simeq 0.235 n_{e0,SEP} T_{e0,SEP}^{\frac{3}{2}}$$

- \succ The peak divertor heat flux induced by ELM is proportional to n_{e0} SEP T_{e0} SEP $^{3/2}$
- > The free-streaming expression is used in the flux-limited expression of the parallel thermal conduction (T.Y. Xia et al., NF 2013)



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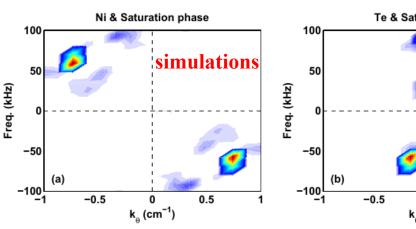


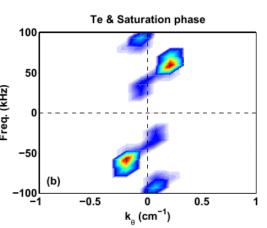


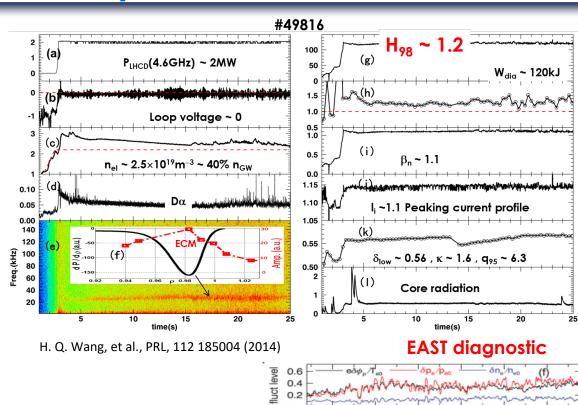
The frequency and poloidal wavenumber of the coherent mode at EAST pedestal are reproduced by the simulations

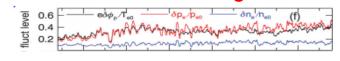
Validations on coherent mode:

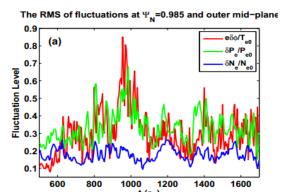
- Similar calculation of f and k
- Simulated frequency is close to exp.
- $k_{\rm e}$ is in the range of exp. (0.5-0.7cm⁻¹)
- Density fluctuations rotate in the electron diamagnetic direction
- The simulated fluctuations are in the similar amplitude with exp.

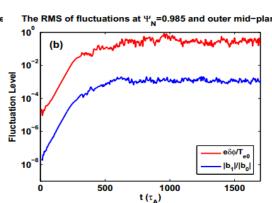








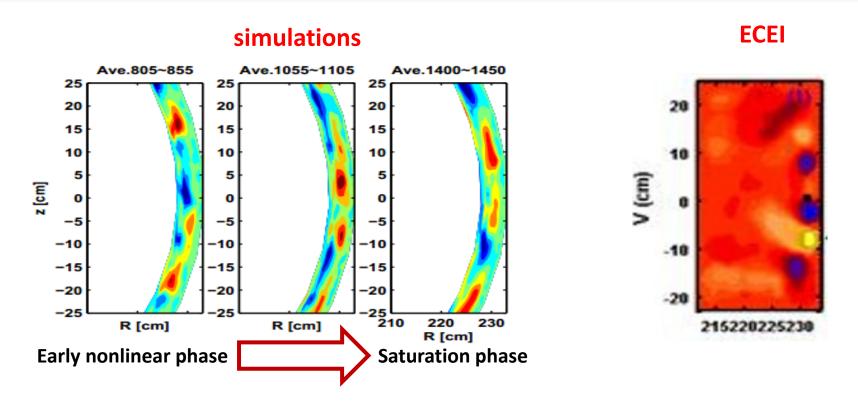








The simulated poloidal mode structures qualitatively agree with ECEI diagnostic



- > The simulated toroidal mode structure shows the similar numbers
- > The radial extension is narrower due to the noise of diagnostic (M. Kim, NF 2014)
- > The poloidal extension is longer, smaller Te wavenumber in simulation.



The mode coupling leads to the generation of the coherent mode

Bicoherence:

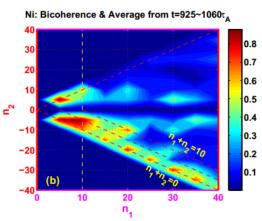
$$S(n_1, n_2) = Y^*(n_3)Y(n_1)Y(n_2), \ n_3 = n_1 + n_2.$$

$$\frac{b(n_1, n_2)}{\sqrt{|Y(n_3)|^2} \sqrt{|Y(n_1)Y(n_2)|^2}}$$

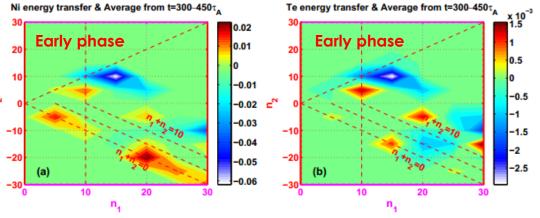
$$\frac{\partial \varphi(k,t)}{\partial t} = \Lambda_{\mathbf{k}}^{\mathbf{L}} \varphi(k,t) + \sum_{k_1 \geq k_2, k=k_1+k_2} \Lambda_{\mathbf{k}}^{\mathbf{Q}}(k_1,k_2) \varphi(k_1,t) \varphi(k_2,t)$$

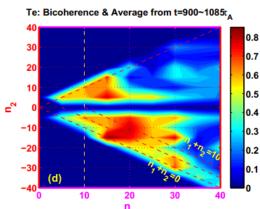
Energy transfer rate*:

$$\frac{\partial P_k}{\partial t} \approx \gamma_k P_k + T_k, \qquad P_k = <\phi_k \phi_k^*>$$



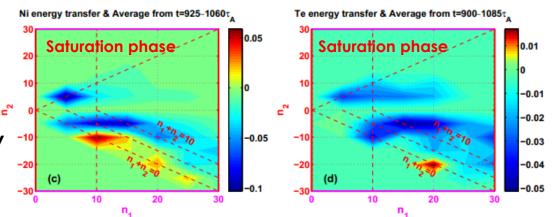
Early nonlinear phase:
 Ni: more modes
 coupled to n=0 then n=15
 Te: more modes
 coupled to n=15 then n=0





Saturation phase:
 Ni: lower n modes
 coupled to n=0
 Te: n=0 mode mos

Te: n=0 mode mostly from (20,-20)





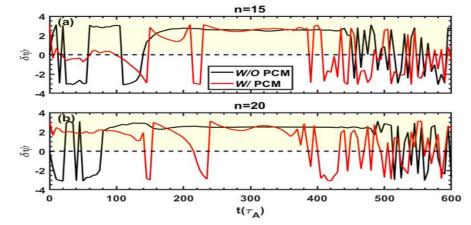


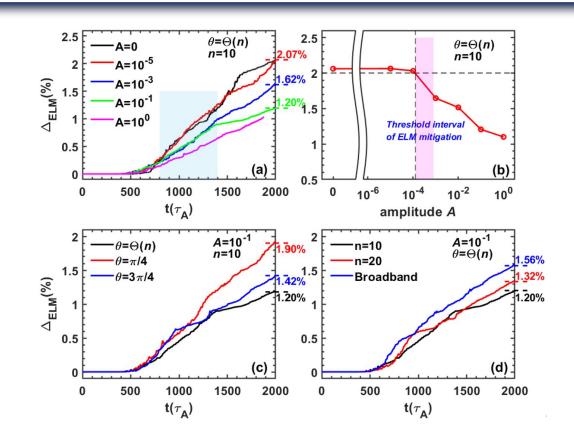
The coherent modes at pedestal are able to mitigate ELM

A modeled pedestal coherent mode (PCM) is added in BOUT++ to simulate the interactions with ELM.

$$\begin{split} \tilde{P}_{PCM}(x,y,z,A) &= Af(z) \cdot e^{\frac{(x-b_1)^2}{2\sigma_x^2} \frac{(y-b_2)^2}{2\sigma_y^2}}. \\ f(z) &= [F(n,\theta)]_{IFFT}. \\ \\ \tilde{P}_{PCM}(x,y,z,A) &= A \cdot e^{\frac{(x-b_1)^2}{2\sigma_x^2} \frac{(y-b_2)^2}{2\sigma_y^2}} [P_{kz}(n)e^{-i\theta}]_{IFFT}. \end{split}$$

- A threshold of the amplitude of PCM is found for ELM mitigation
- Phase angle and dominant toroidal mode can also affect ELM.





The phase coherent time* is changed by the nonlinear wave-wave interactions in the linear phase, which change the growth time of the ELM.



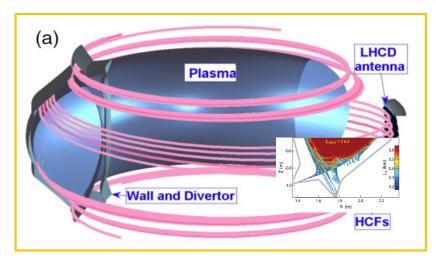
Y.L. Li et al., submitted to NF *P.W. Xi, X.Q. Xu and P.H. Diamond, 2014 PRL 112



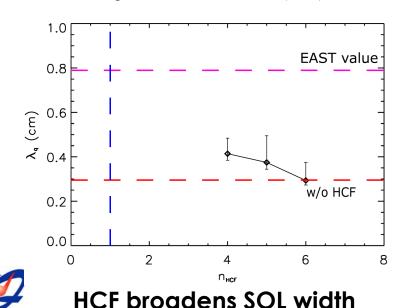
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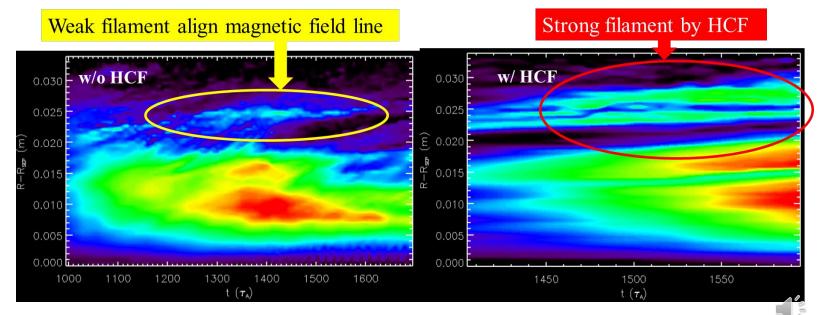
The simulations reproduce the splitting of the strike point by helical filamentary current (HCF)



Y.F. Liang et al., PRL 110, 235002 (2013).

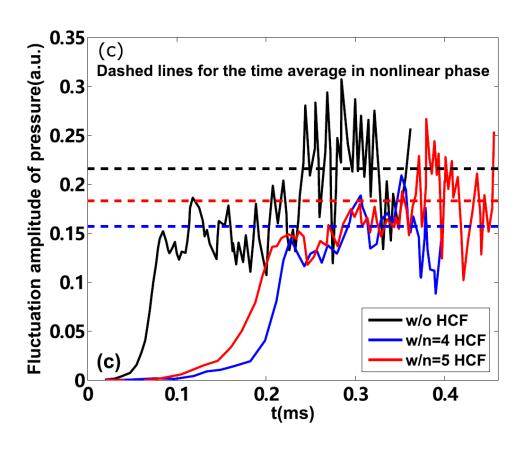


- > A modeled HCF with force-free form in SOL is added into BOUT++ as the extra magnetic flutter.
- ➤ HCFs dominated by n=1, but n>1 used in the simulation due to the efficiency
- \triangleright SOL width λ_a is indeed broadened by HCFs
- > The splitting of strike point behavior is reproduced



T.Y. Xia et al., NF 59 076043 (2019)

The simulation proves that LHW can mitigate the edge turbulence effectively through the HCFs



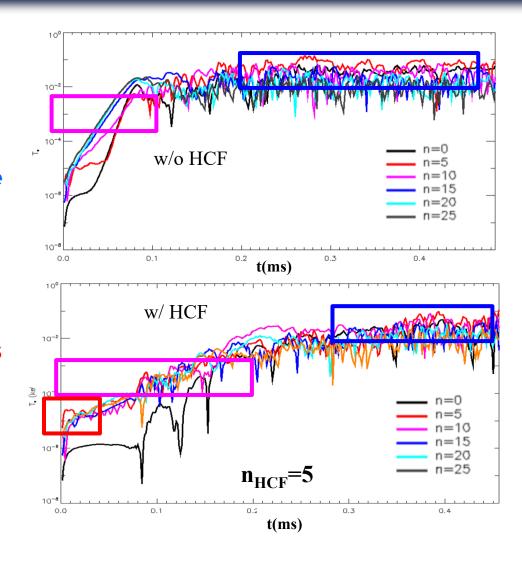
- \succ The amplitude of fluctuation is decreased by HCF \rightarrow mitigation of ELM
- > HCF with lower n shows more obvious mitigation on fluctuations





The HCF induced Br decreases the linear growing of the fluctuations

- > At linear phase, the dominant mode w/ HCF is decreased towards the number of HCF
- > At late nonlinear phase, the dominant modes are the same.
- > The linear growing is mitigated by Br induced by HCF.
- ➤ Nonlinear mode interactions among multi modes driven by HCF Br decrease the phase coherent time (P.W. Xi, PRL, 2014) and leads to the slow growing.



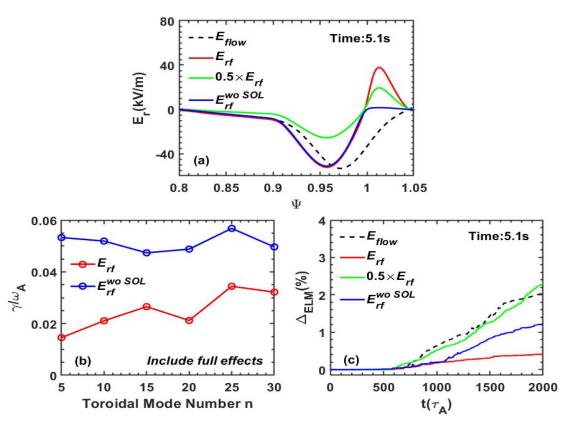


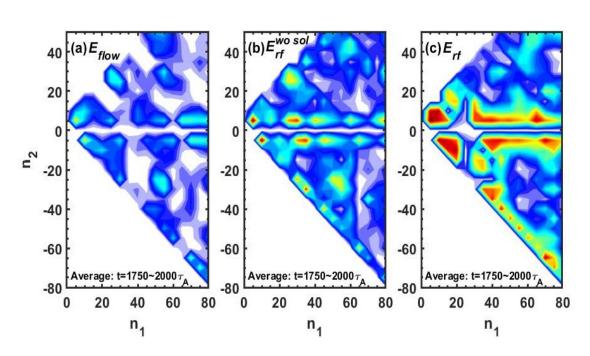
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The RF sheath by ICRF may suppress ELM effectively

- > The RF sheath is able to suppress ELM from linear phase.
- > The RF sheath generates the strong flow shear near the separatrix*, which leads to the widely nonlinear mode coupling.





more nonlinear modes coupling with the effects of the sheath potential





Summary

- > The 6-field 2-fluid model in BOUT++ framework is developed to study the edge turbulence and ELMs for the typical RF heating H-mode on EAST
- > The RF effects can be studied from 2 aspects:
- ➤ On equilibrium:
 - The pedestal scan shows the peak value of the divertor heat flux by ELM is proportional to $n_{e0\ SEP}$ $T_{e0\ SEP}^{3/2}$
 - The mode coupling process generates the edge coherent mode. This mode can mitigate ELM.
- > On instabilities:
 - HCF by LHW is able to mitigate the edge instabilities and broaden the SOL width
 - RF sheath by ICRF is able to suppress ELMs

More integrated simulations are necessary!





